

USPAS - Fundamentals of Ion Sources

7. Multicusp Ion Sources I

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Damon Todd (LBNL),
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Outline

- Introduction – basic principles of multicusp ion sources
- Different heating methods (filament, RF antenna, ECR?)
- Usages: H^+ , H^- , He^+ , H_2^+ , ...
- Positive (H^+ , H_2^+ , He^+ , ...) multicusp ion sources
 - Optimization
- Negative (H^- , ...) multicusp ion sources
 - Volume production
 - Surface-enhanced volume production (Cs)
 - Extraction of negative ions
 - Optimization
- Other types of negative ion sources

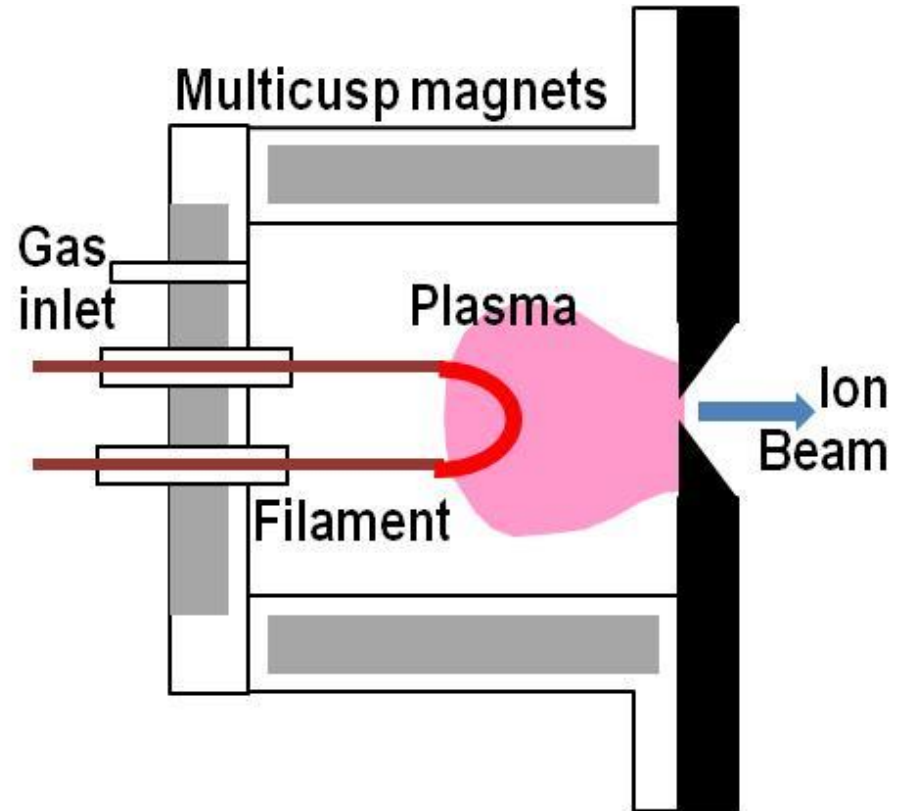
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Basics

Basic Building Blocks:

- Confinement
- Raw Material (Gas)
- Energy for Electrons (Heating)
- Extraction
- 1973 Limpaecher & MacKenzie at UCLA build a 86 liter plasma vessel with 1252 alternating Alnico bar magnets lining the walls. Plasma is very quiescent; RSI 44, 1973, 726
- 1979 Ehlers and Leung at LBNL start to develop multicusp ion sources for hydrogen and later for negative ions.



Design Parameters

Parameter	Typical Values	Tuning	Comment
Magnetic Field	?	?	
Heating	?	?	
Material	?	?	
Extraction	?	?	

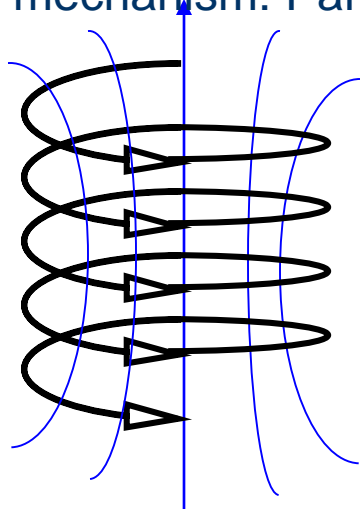
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Heating	?	?	
Material/Gas	?	?	
Extraction	?	?	

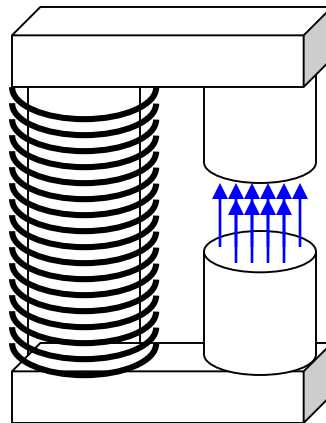
Start with magnetic field...

Magnetic Confinement

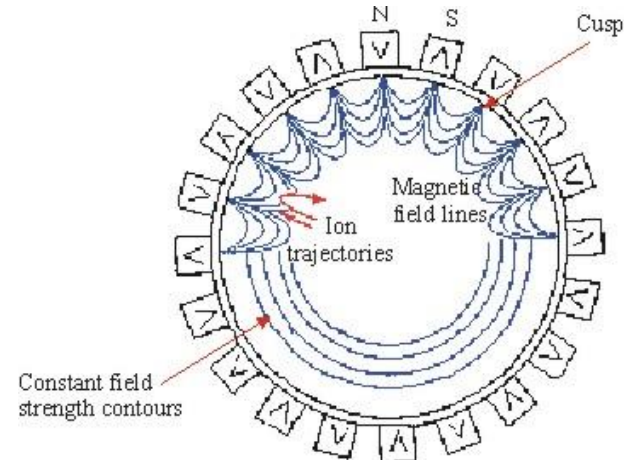
- In the direction of the fields there is no force, transverse the particles are bend into the circular motion $F = q(\vec{E} + \vec{v} \times \vec{B})$
- Helical motion increases the time the electron send in the discharge chamber- field lines can only be crossed through collisions – wall losses are reduced
- Add strongly increasing magnetic field as the confinement mechanism: Particles get reflected by an increasing magnetic field



Solenoidal field

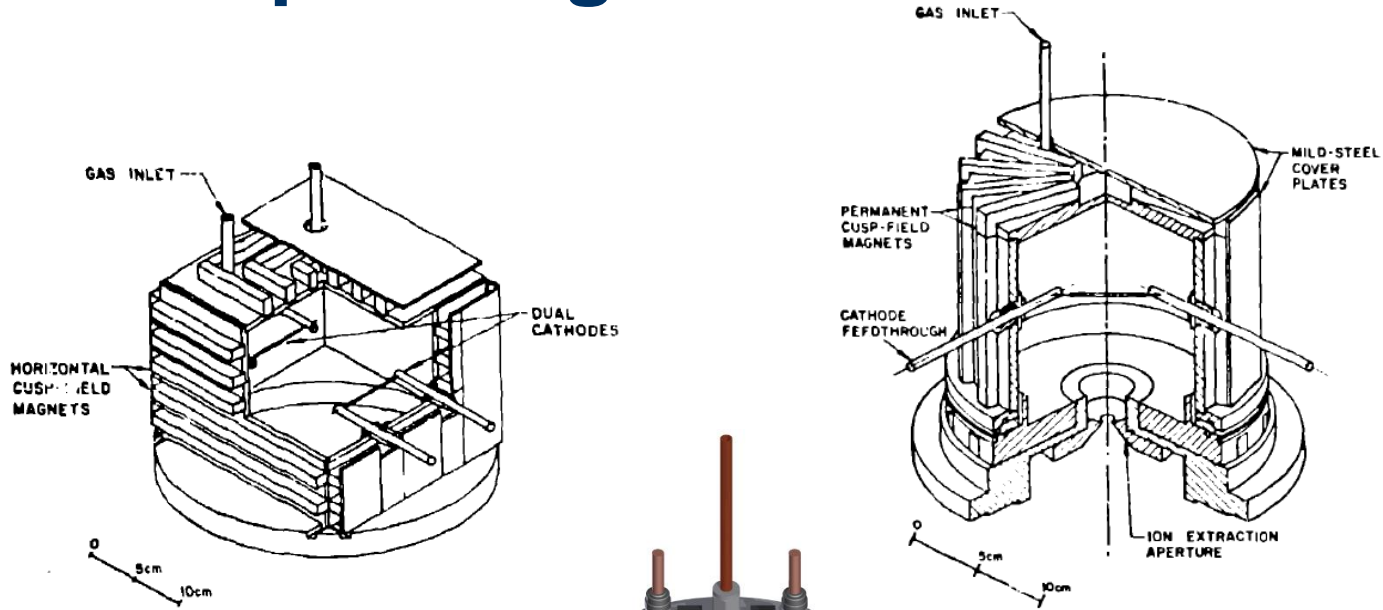


Dipole field

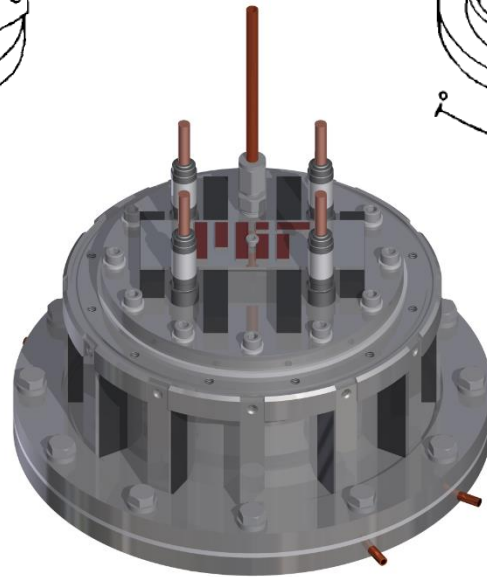


Multicusp field

Multicusp Configurations

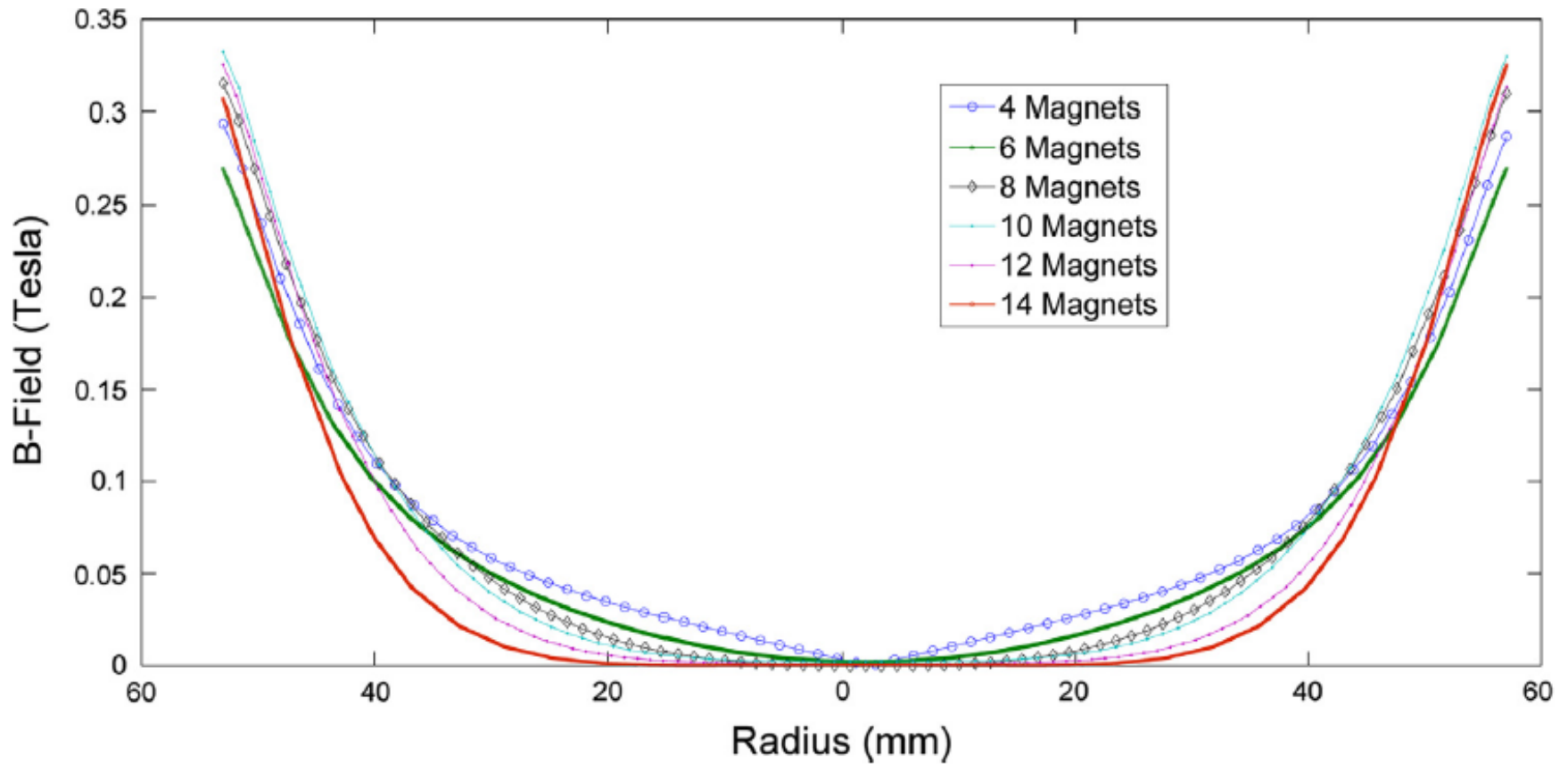


D. Boonyawan et al., RSI 71, 2000

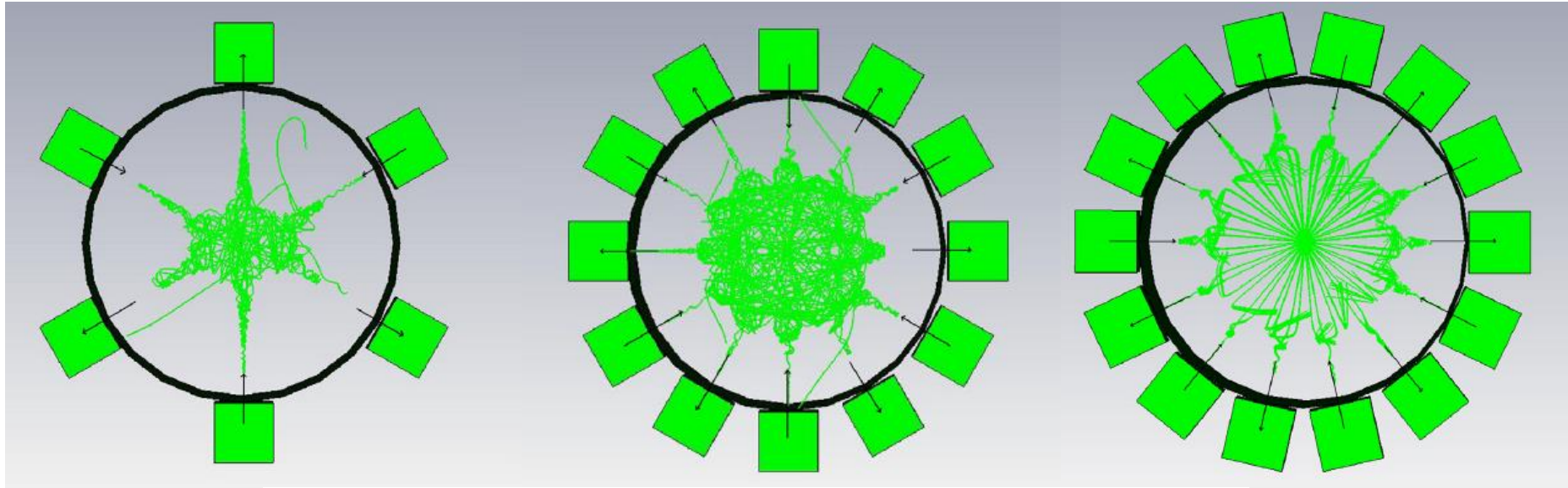


S. Axani, D. Winklehner, MIT 2015-2016

Number of Magnets I



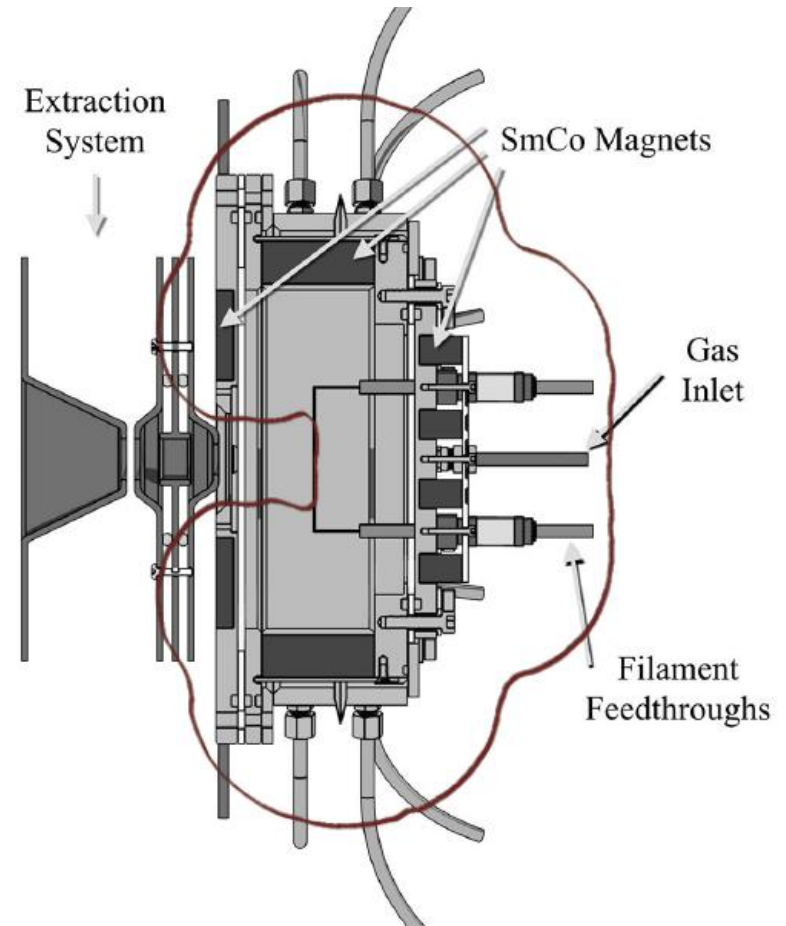
Number of Magnets II



Mean life of electrons(s)	Sum of electron trajectory (m)	Number of magnets
$7.06e-7$	29,453	4
$1.47e-6$	61,641	6
$5.51e-6$	231,440	8
$7.8e-6$	328,062	10
$8.79e-6$	368,032	12
$7.01e-6$	294,260	14

Field Free at Extraction

- We typically use the 10 Gauss contour to define a field-free region (< 10 Gauss).
- Arrange magnets such that there is a field-free region around extraction.
- Should be able to achieve very low emittances. Now only depends on T_i , extraction system and I_{beam}



Permanent Magnets

Magnet	B_r (T)	H_{ci} (kA/m)	$(BH)_{max}$ (kJ/m ³)	T_c (°C)
Nd ₂ Fe ₁₄ B (sintered)	1.0–1.4	750–2000	200–440	310–400
Nd ₂ Fe ₁₄ B (bonded)	0.6–0.7	600–1200	60–100	310–400
SmCo ₅ (sintered)	0.8–1.1	600–2000	120–200	720
Sm(Co,Fe,Cu,Zr) ₇ (sintered)	0.9–1.15	450–1300	150–240	800

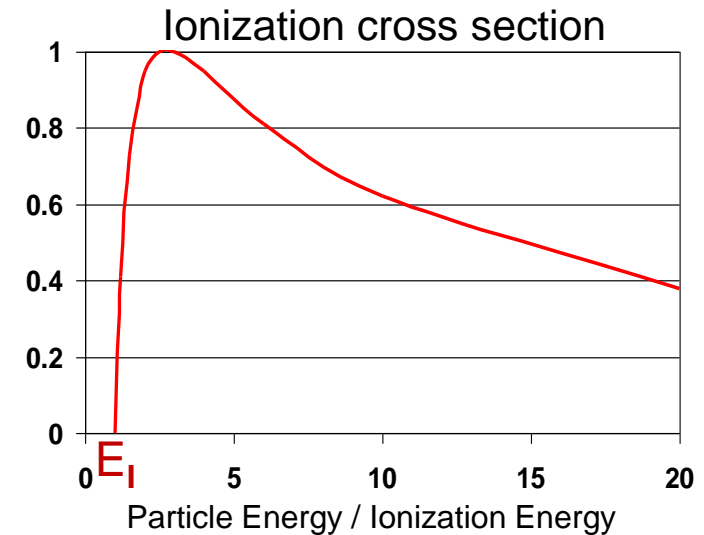
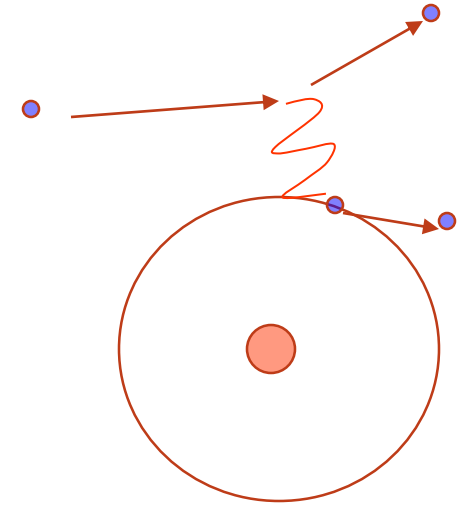
Source: Wikipedia

Design Parameters

Parameter	Typical Values	Tuning	Comment
Permanent Magnet Multicusp Magnetic Field	0.8 – 1.4 T 8 – 16-Pole Front/Sides/ Back	No	Most likely need cooling NdFe has higher field SmCo better Tc
Heating	?	?	
Material/Gas	?	?	
Extraction	?	?	

Electron Impact Ionization

- The removal of an electron from gaseous atoms or molecules requires electric fields in excess of 10^{10} V/m, only possible with atomic distances typically reached in collisions with charged particles.
- The conservation of energy and momentum favors electrons as the most efficient ionizing particles, and therefore most ion sources use electron impact ionization.
- The conservation of energy is responsible for an absolute threshold, the ionization energy E_i , the minimum energy which needs to be transferred for successful ionization.
- Gases have ionization energies between 12 eV for O_2 and 25 eV for He, e.g. 15 eV for H_2 molecules and 14 eV for H atoms.
- The ionization cross section has a maximum close to 3 times the ionization energy E_i and therefore electrons with an energy between 50 and 100 eV ionize all gases efficiently.



Heating Methods

How can we give the electrons the necessary energy to ionize?

- Filament: Create through heat and accelerate in low static electric field.
- RF: Use electrons from collisions and accelerate with electric field
- Electron Cyclotron Resonance (ECR). 2.45 GHz corresponds to ~ 875 Gauss.

Townsend Discharge

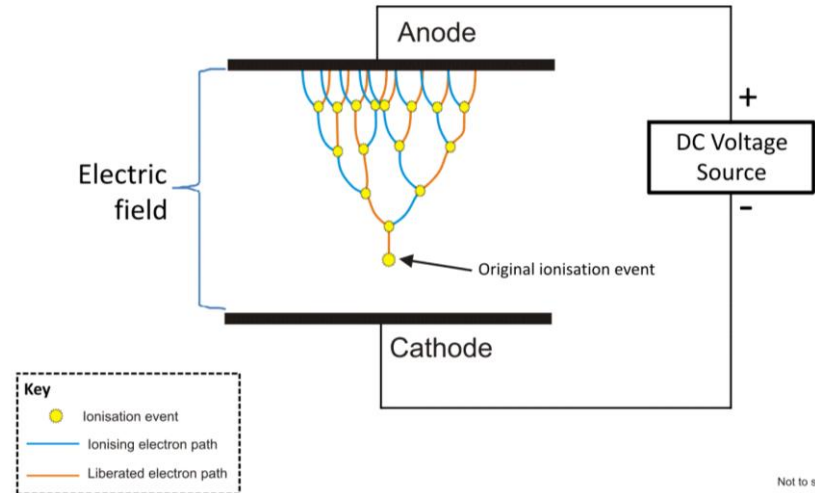
- Around 1900 John Sealy Townsend studied discharges below the breakdown voltage (dark discharges).
- UV light illuminating the cathode produces photo electrons.
- Applying sufficient voltage, the discharge current I grows exponentially with the distance d between the electrodes:

$$I = I_0 \cdot e^{\alpha \cdot d}$$

where I_0 is the photoelectric current and α is the 1st Townsend coefficient

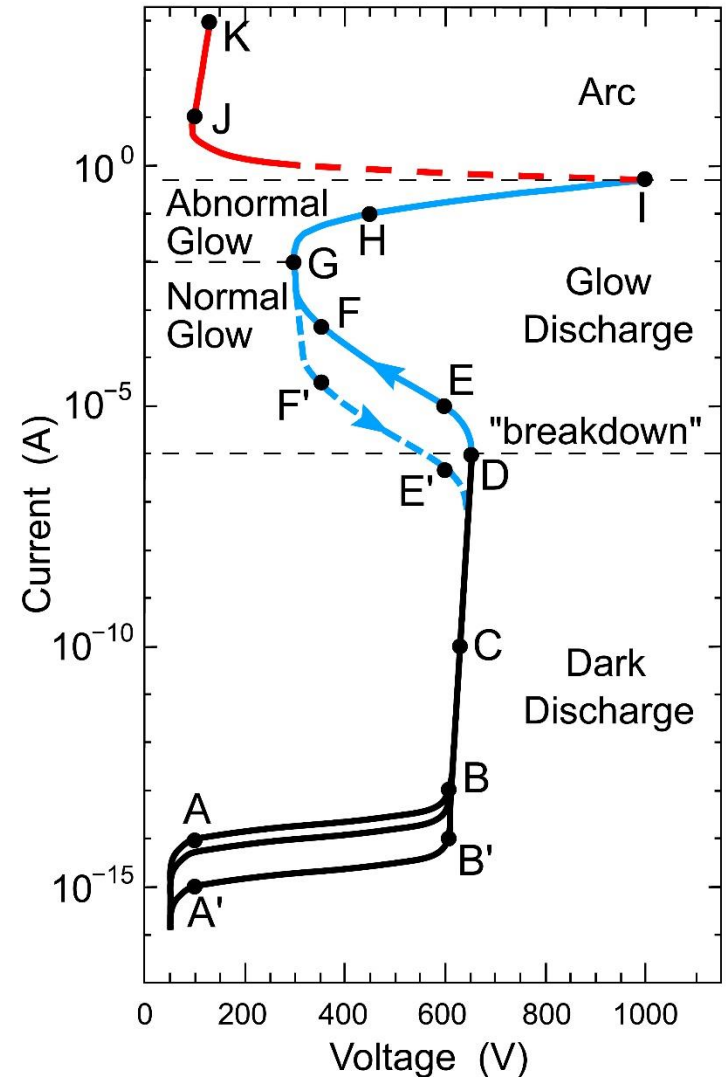
- Increasing the gap d increases the number of electron multiplications, causing an avalanche, i.e. exponential growth.
- Keeping d constant and increasing the voltage V shortens the distance between ionizing collisions, which increases the discharge current I exponentially.

Visualisation of a Townsend Avalanche



Discharge Regimes

- Small voltages yield nA currents by collecting electron–ion pairs produced by radiation.
 - Raising the voltage starts the Townsend multiplication, yielding many μA (corona).
 - Suddenly the **gas starts to glow** and the current grows to many mA at a much reduced voltage.
 - Discharge current diverges when the ions impacting on the cathode generate enough secondary electrons to replace all seed electrons:
- $$I = I_0 \cdot \frac{e^{\alpha \cdot d}}{1 - \gamma \cdot (e^{\alpha \cdot d} - 1)}$$
- Increasing the reduced voltage increases the **glowing plasma volume** and the current.

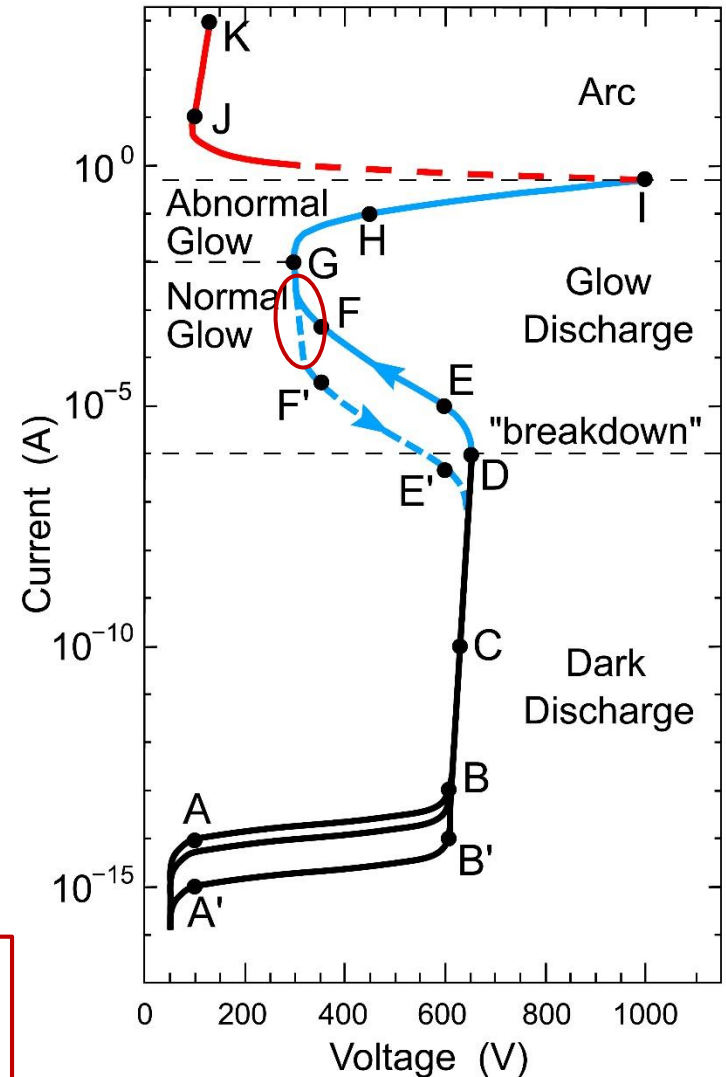


(Adapted from M. Stöckli, ORNL)

Discharge Regimes

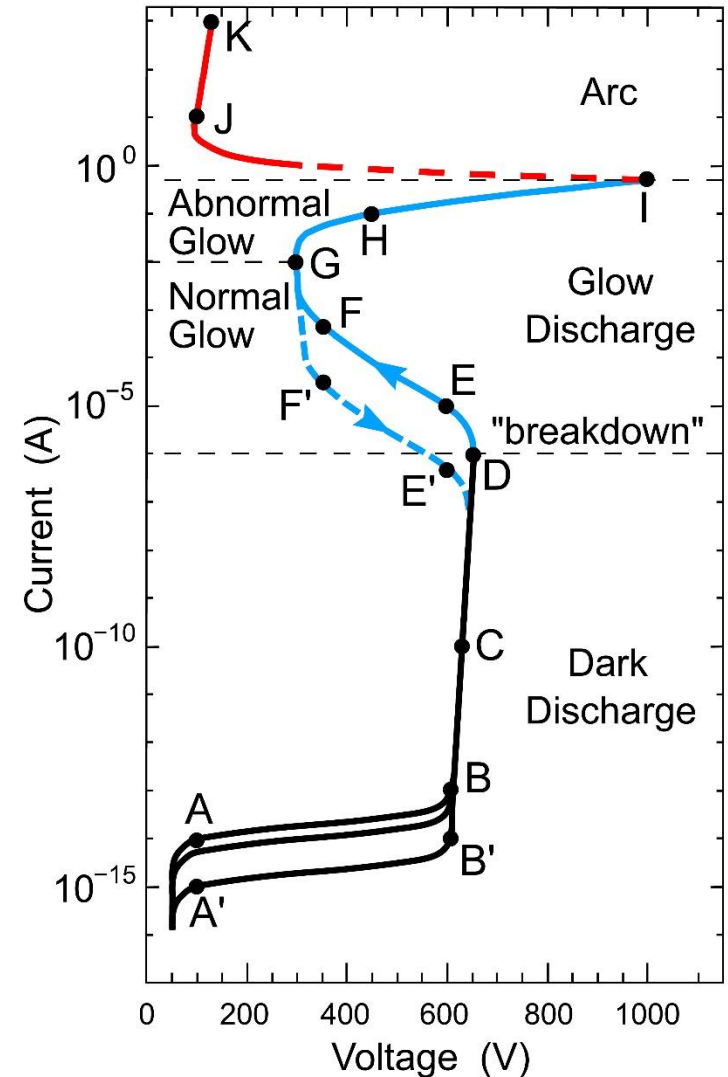
- Small voltages yield nA currents by collecting electron–ion pairs produced by radiation.
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Discharge ion sources typically operate at the low current end of glow discharges.



Discharge Regimes

- A. random pulses by cosmic radiation
- B. saturation current
- C. avalanche Townsend discharge
- D. self-sustained Townsend discharge
- E. unstable region: corona discharge
- F. sub-normal glow discharge
- G. normal glow discharge
- H. abnormal glow discharge
- I. unstable region: glow-arc transition
- J. electric arc
- K. electric arc



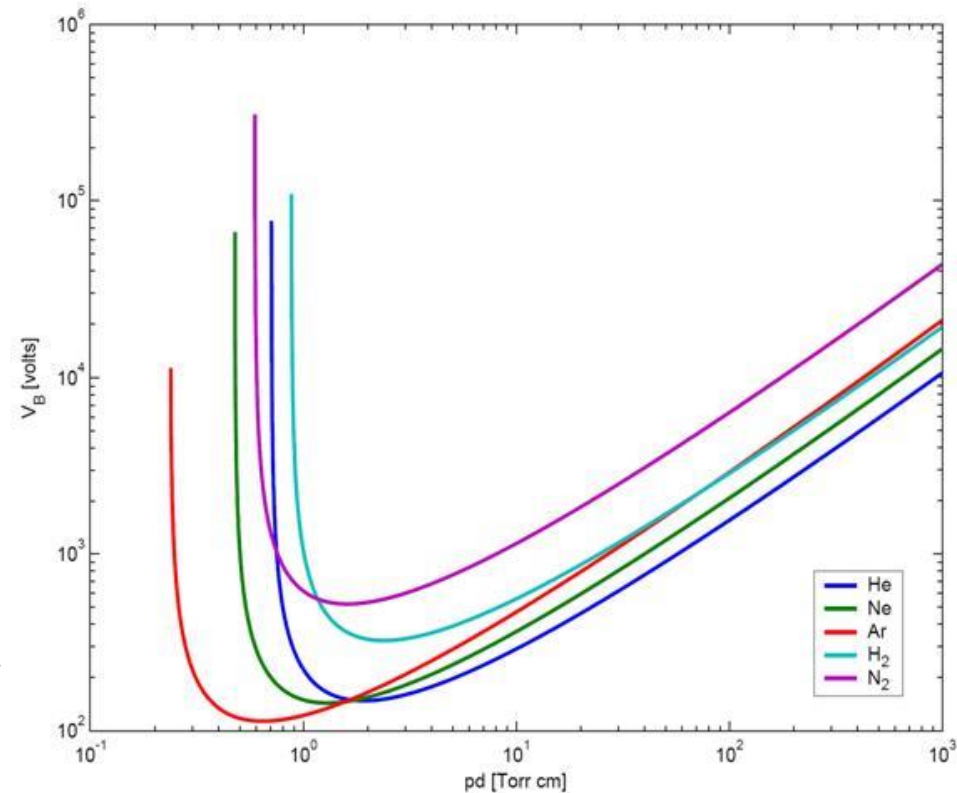
Breakdown Voltage (Paschen's Law)

- In 1889, Friedrich Paschen describes the breakdown voltage V :

$$V = \frac{a \cdot p \cdot d}{\ln(p \cdot d) + b}$$

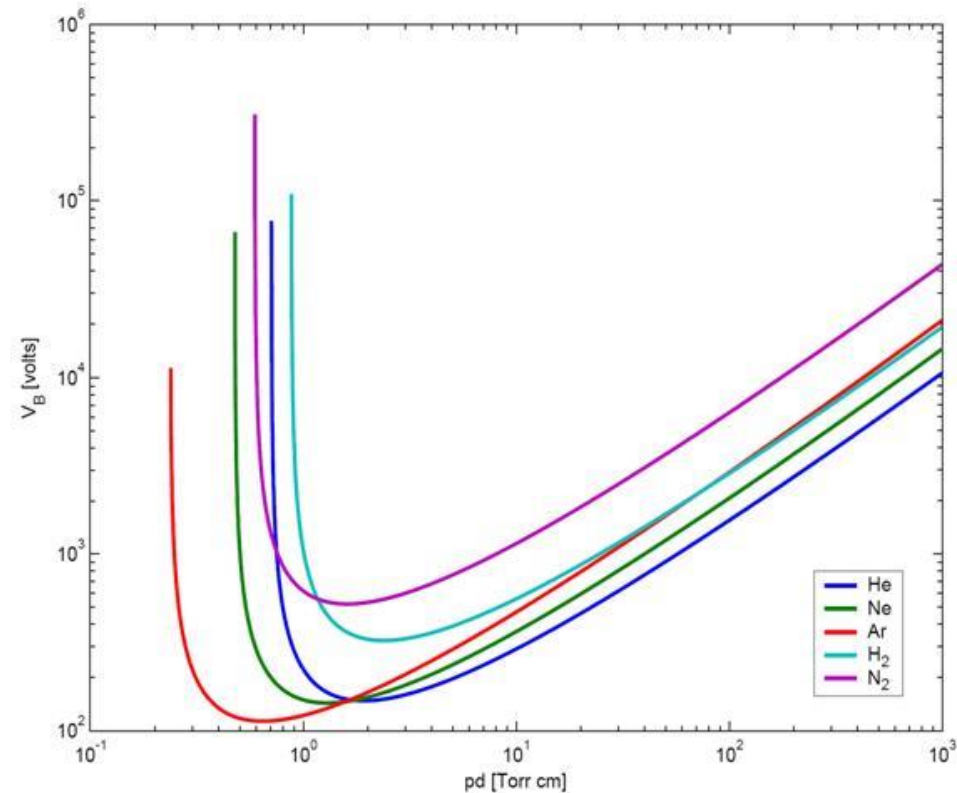
with pressure p , electrode gap d , and parameters a & b , which depend on the gas and the electrodes.

- Normalization with the minimum voltage V_{min} and corresponding $(p \cdot d)_{min}$ creates a universal curve.
- Gases are good insulators at very low and very high pressures.



Breakdown Voltage (Paschen's Law)

- Free electrons have to gain enough energy to ionize an atom in the next collision (λ_i).
- Decreasing the pressure increases the mean path between collisions (λ_i), which is compensated by proportionally increasing d .
- The minimum represents the minimum energy spent on producing enough ions for one secondary electron from the cathode.
- At high $p \cdot d$, the voltage increases linearly with the gap between the electrodes.
- At low $p \cdot d$, the electrons need more energy to liberate more than one electron.



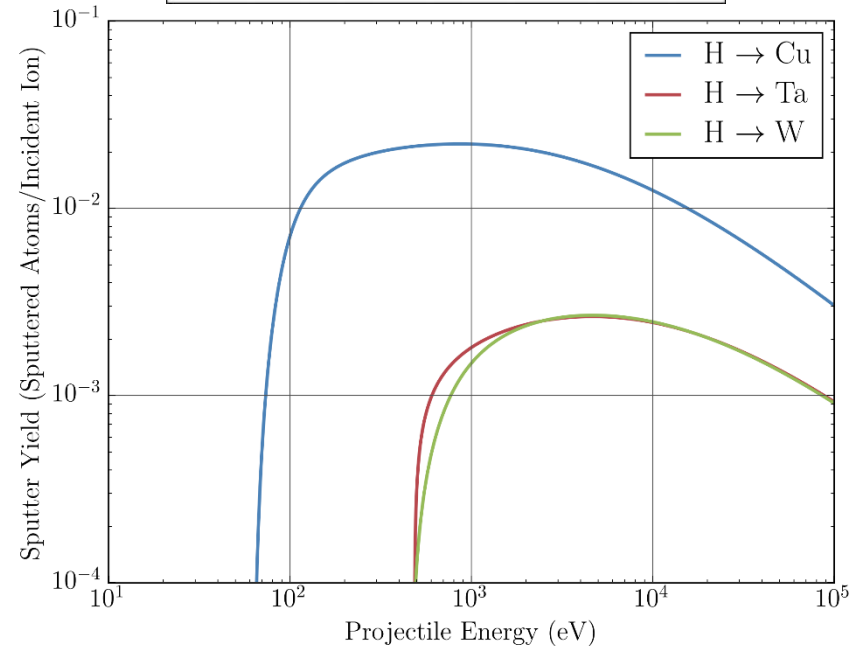
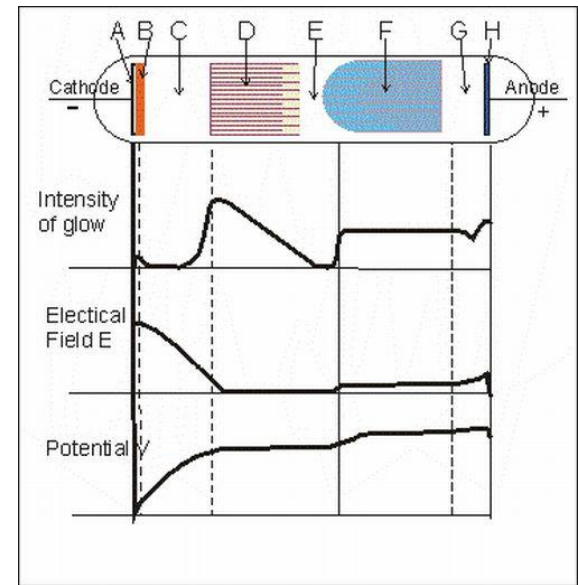
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Permanent Magnet Multicusp Magnetic Field	0.8 – 1.4 T 8 – 16-Pole Front/Sides/ Back	No	Most likely need cooling. NdFe has higher field SmCo better Tc
<u>Filament Heating</u>			
Heating Current	O(10A)	Yes	
Discharge Current	O(1A)	Yes	
Discharge Voltage	O(100V)	Yes	
Filament Position	O(10cm)	Maybe	Depends on Ion
Material/Gas	1e-5 to 1e-3 Torr	Yes	
Extraction	?	?	

Filament Lifetime?

- An electric field is required to accelerate the electrons to an energy sufficient to ionize the neutral particles.
- The electric field, however, also accelerates the ions in the plasma sheath at the cathode. These accelerated ions impact on the electrode
- Sputtering reduces the filaments thickness until they break.
- Sputtered metal atoms coat insulators until they break down.

Sputtering reduces the lifetime of the source and needs to be minimized!



Improve the Situation

- The low sputter rates of heavy, refractory metals yield longer lifetimes!

Material	Sputter Rate @ 1 keV (atoms/ion)	Sputter Threshold (eV)
Cu	0.03	O(60)
W	0.0015	O(450)
Ta	0.0018	O(450)

LaB₆ Cathodes:



- Insulator lifetimes can be extended with recessed areas providing partial shadows. This extends the lifetime until the growing metal films flake and peel away from non-shadowed areas and short out an insulator. (not only a problem with filaments!)
- What if we don't need a filament? → RF antennas (but come with other problems)

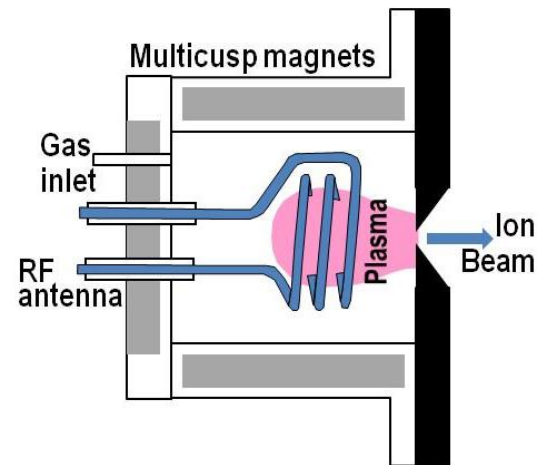
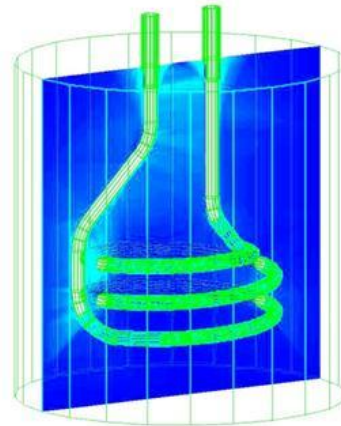
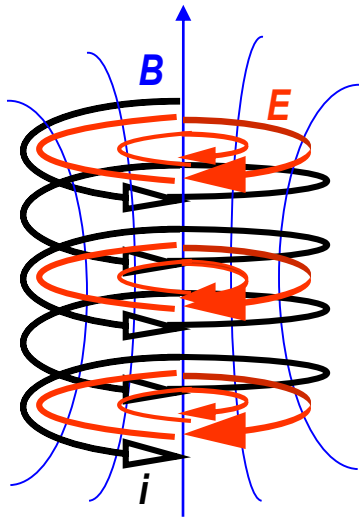
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Discharge Current	O(1A)	Yes	
Discharge Voltage	O(100V)	Yes	
Filament Position	O(10cm)	No-ish	Depends on Ion
Filament Material	W, Ta, LaB ₆	No-ish	Filament Shape? Lifetime
Material/Gas	1e-5 to 1e-3 Torr	Yes	
Extraction	?	?	

RF Driven Multicusp Ion Sources

- The 2nd Maxwell Equation describes a curling E field generated by a changing magnetic field in absence of any charges! $\Delta \times E = -\frac{\partial B}{\partial t}$
- A changing magnetic field B can be produced with an alternating current $i = i_0 \cdot \cos(\omega t)$ in N windings with radius r_0 : $B = \frac{1}{2}\mu_0 \frac{N \cdot i}{r_0}$
- Now integrate Maxwell's equation to get Faraday's law and solve for E :

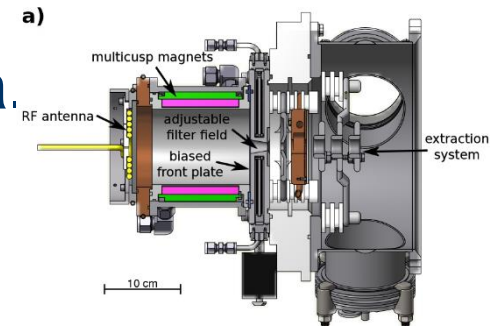
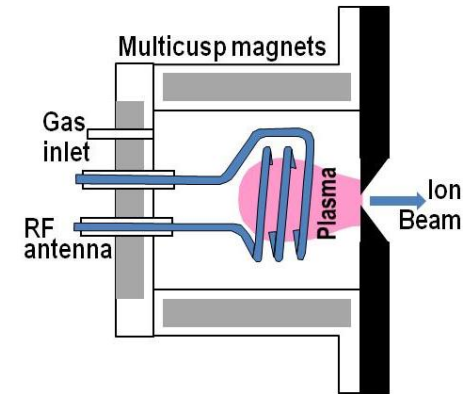
$$\int E \cdot ds = -\frac{d\Phi_B}{dt} = -\frac{d}{dt} \int B \cdot dS \rightarrow E(r, t) = \frac{1}{4} \frac{r}{r_0} \cdot \mu_0 \omega N i_0 \cdot \sin(\omega t)$$



Adapted from: M. Stöckli, US-CERN-Japan Accelerator School 2013

Types of RF drives

- Capacitively coupled (not typically used for ion sources)
- Inductively coupled with Internal Antenna.
- Inductively coupled with external Antenna.

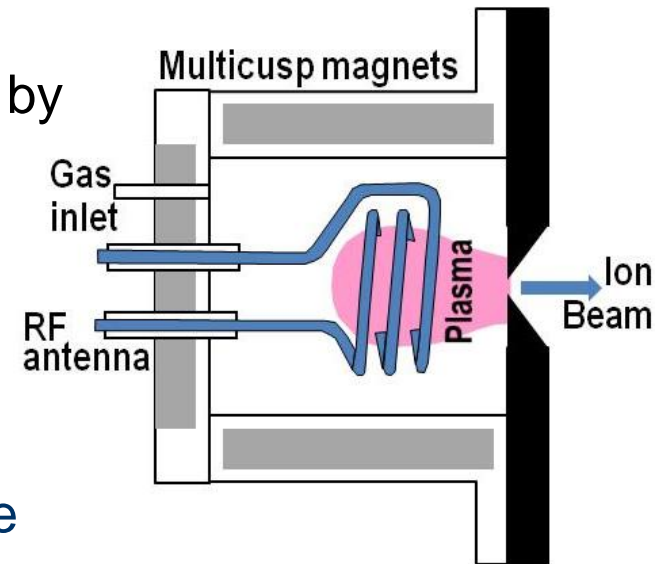


- In 1990 LBNL replaced the filament in their multicusp ion source with a 3-turn antenna, driven by 2 MHz.
- In 1992-1994 this source was tested for SSC yielding up to 100 mA (?) for 0.1 ms at 10 Hz.
- 1999-2001 LBNL developed this source for SNS, hoping to eventually reach the 7% duty factor.

In 2001 the plasma was visualized by replacing the stainless steel source chamber with a glass dome surrounded by identical cusp magnets.

- The **inductively induced plasma** is bright **inside** the antenna, driven by the high induced fields near the antenna. The magnetic bucket causes the plasma to drift to the center, compensating for the smaller electric fields.
- Less plasma drifts into the confining cusp fields.

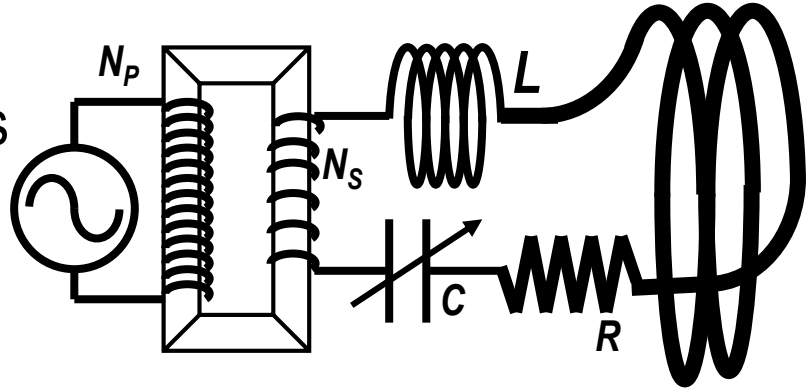
Multicusp confinement is well suited for an inductively driven ion source!



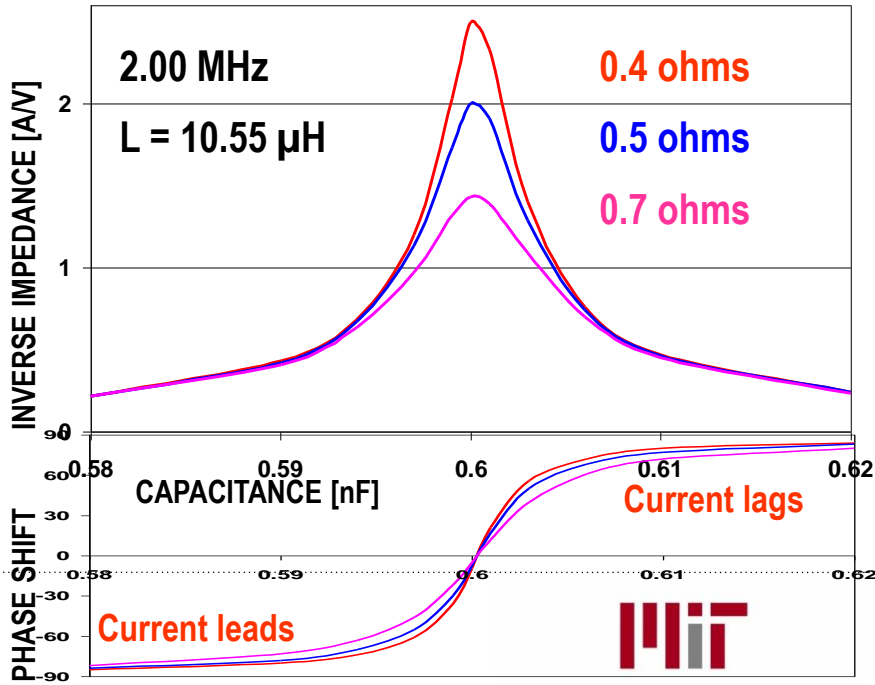
Adapted from: M. Stöckli, US-CERN-Japan Accelerator School 2013

Making more ions with a RF-driven Ion Source

- More ions \Leftarrow denser plasma \Leftarrow higher electric fields \Leftarrow higher antenna current!
- The RF amplifier output impedance needs to be matched to the impedance of the ion source to get the system in resonance!
- The ion source RLC circuit has a resonant frequency of $\omega^2 = (LC)^{-1}$ and an impedance $Z = \epsilon_0 / i_0 = (R^2 + (\omega L - (\omega C)^{-1})^2)^{1/2}$
- $L \approx 10 \mu H$ and $\omega \approx 2 \cdot \pi \cdot 2 \text{ MHz}$ requires tuning C around 0.6 nF to obtain the maximum current $i_0 = \epsilon_0 / R$.
- Adjusting the transformer ratio N_S / N_P matches the RF amplifier power to the impedance of the ion source.
 - The resonance can be located with the phase shift.



ANTENNA CURRENT versus CAPACITANCE

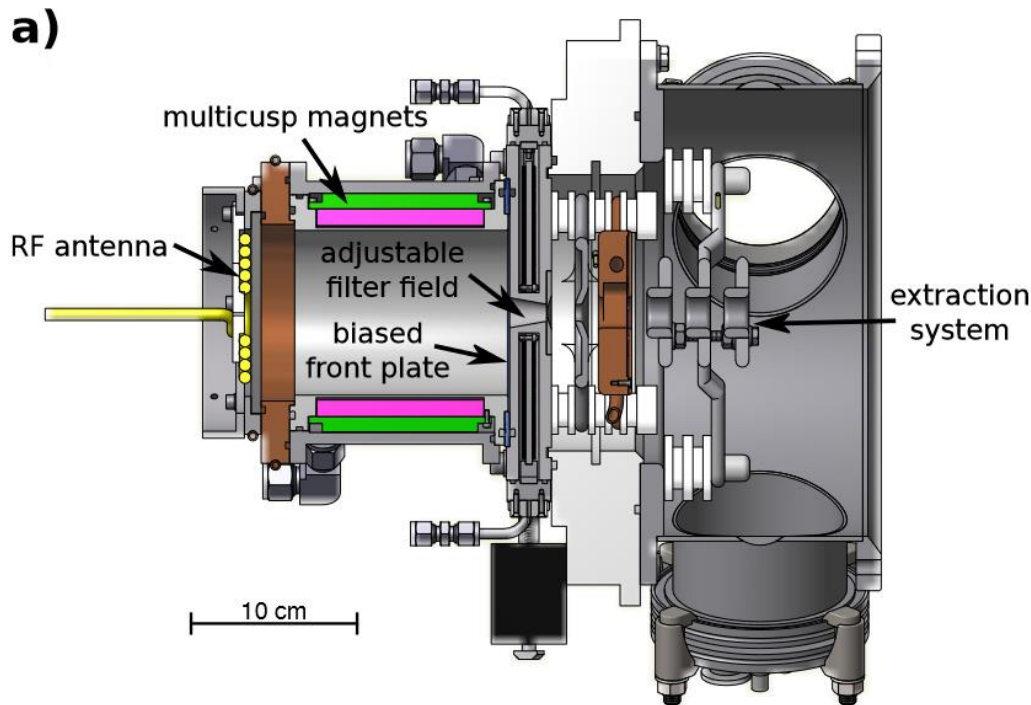


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<u>Filament Heating</u>			Many Parameters
... or ...			
<u>RF Heating</u>			
Frequency	5~13 MHz	Somewhat	
Power	few – 100 kW	Yes	
Material/Gas	1e-5 to 1e-3 Torr	Yes	
Extraction	?	?	

External RF Antennas

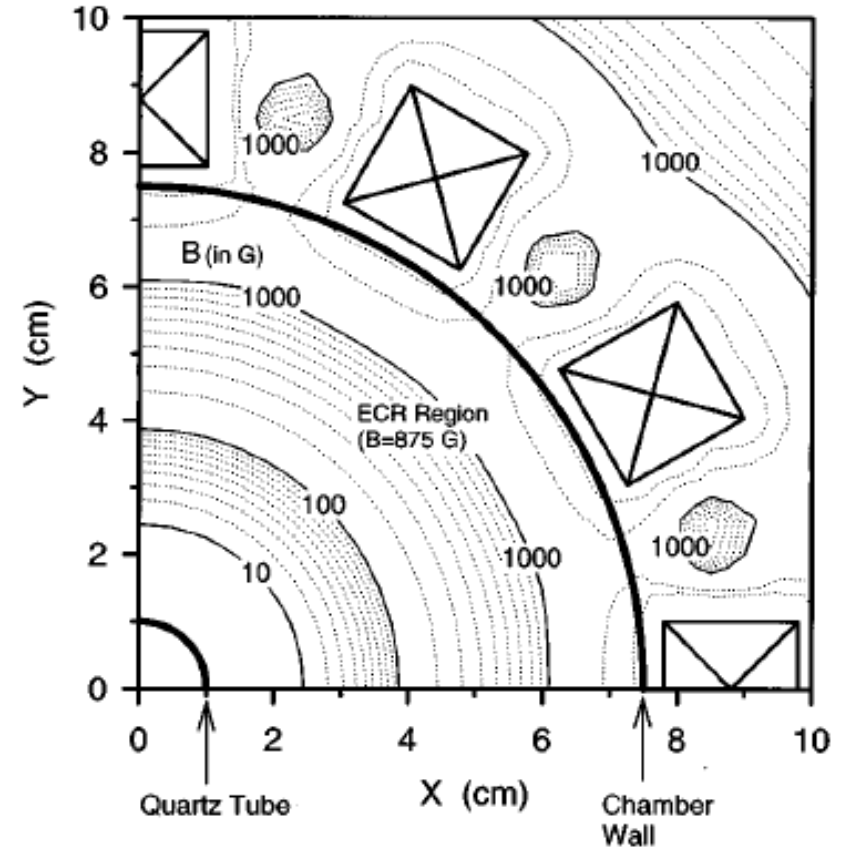
In the back:



Around the chamber:

Hybrid ECR-Multicusp Ion Sources?

- Typical field contours in multicusp source.
- Stands to reason that a 2.45 GHz microwave generator could be coupled through a waveguide (as in an ECR)
- Some promising first results are out there...

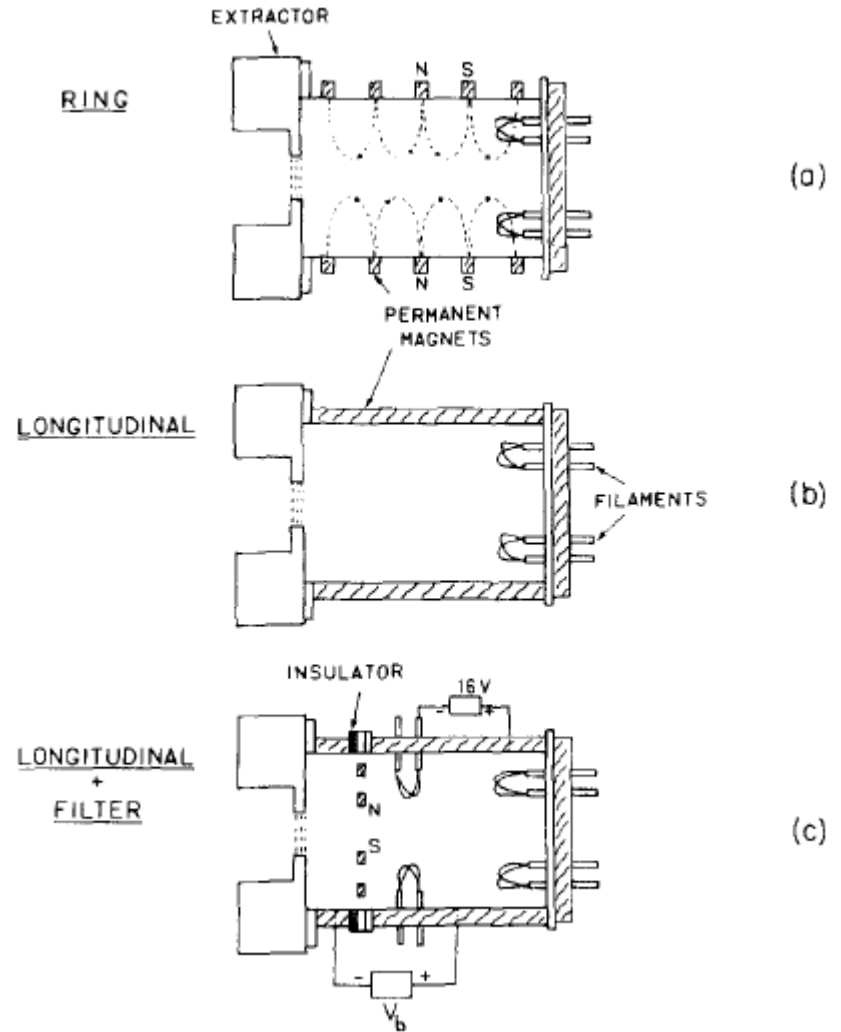
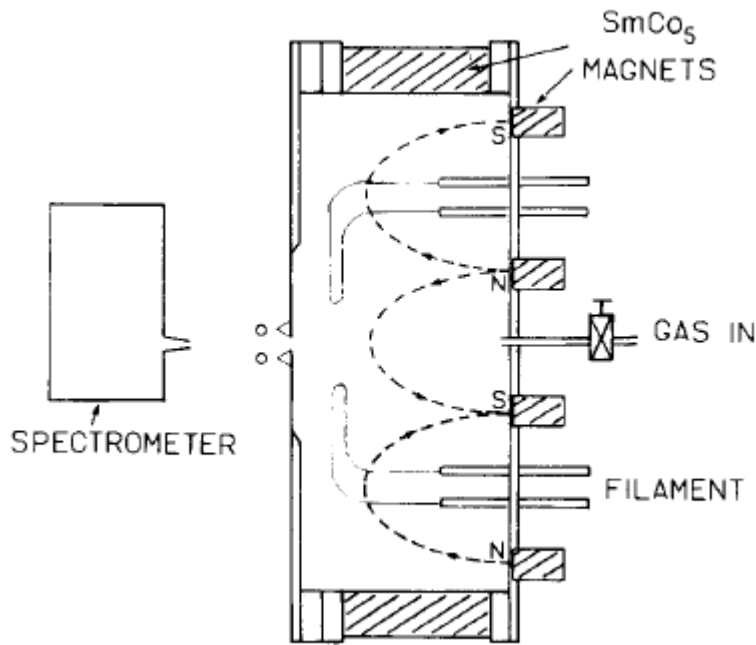


M. Tuda et al., Journal of Vacuum Science and Tech. A, 1998

Examples of Positive Ion Multicusp Sources

LBNL did a lot of development
1970ies – now.

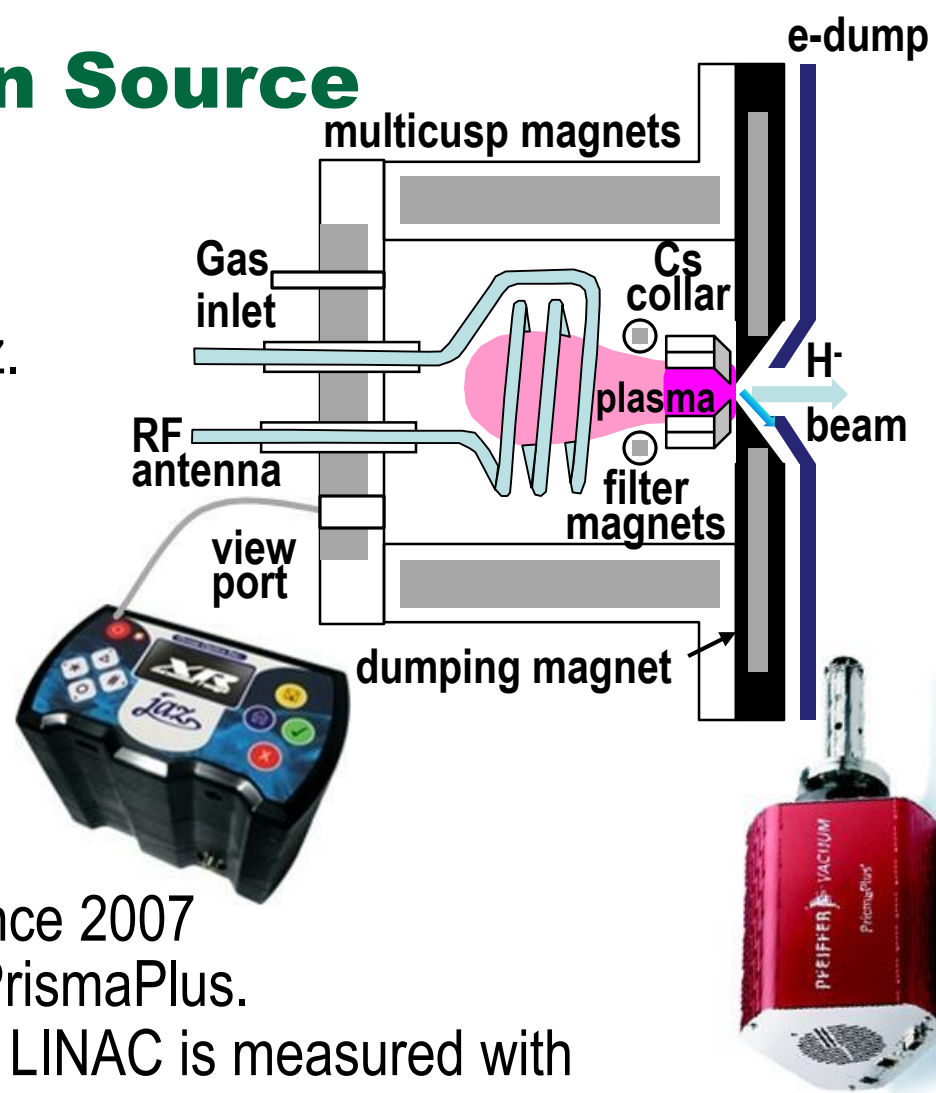
K. Leung, K. Ehlers



The SNS Baseline Ion Source

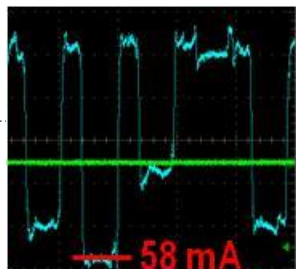
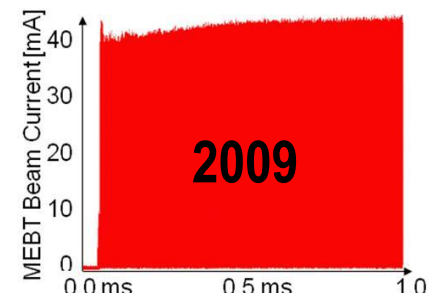
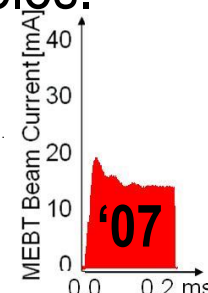
LBNL developed and SNS improved the cesium-enhanced, RF-driven multicusp ion source, which delivers ~1-ms long H^- current pulses at 60 Hz.

- About 300 W 13-MHz RF generate continuous low-power plasma.
- 50-60 kW of 2 MHz RF are added for ~1 ms at 60 Hz to produce the H^- beam pulses.
- A jaz spectrometer analyzes 220-1100 nm emissions since 2011.
- Gaseous emissions are monitored since 2007 and since 2013 with a more sensitive PrismaPlus.
- The H^- beam current injected into the LINAC is measured with a beam current torroid after the RFQ and 2 quadrupoles.

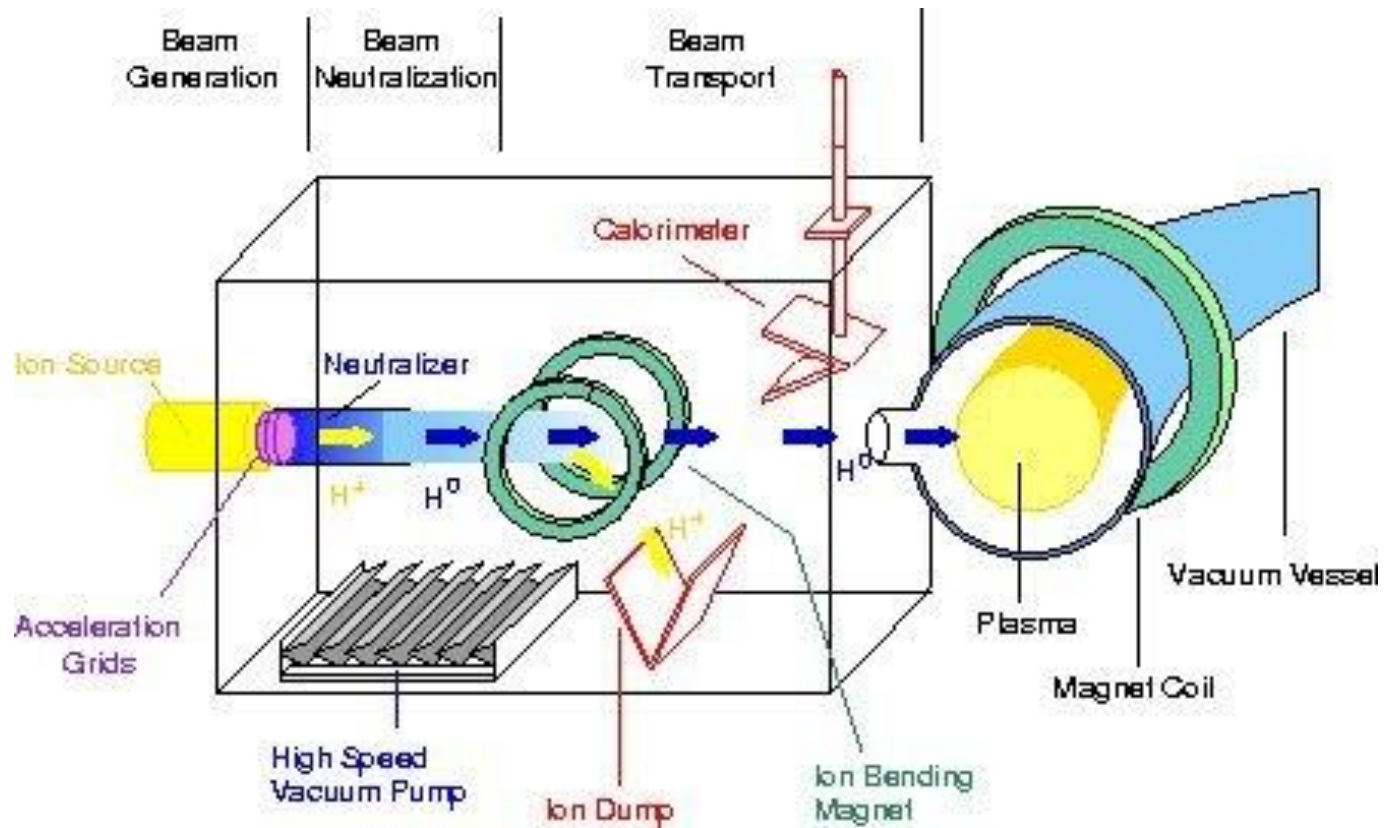


- The beam current injected into the RFQ is measured since 2012.

50-60 mA go into the RFQ!



Giant Ion Sources for Neutral Beams I

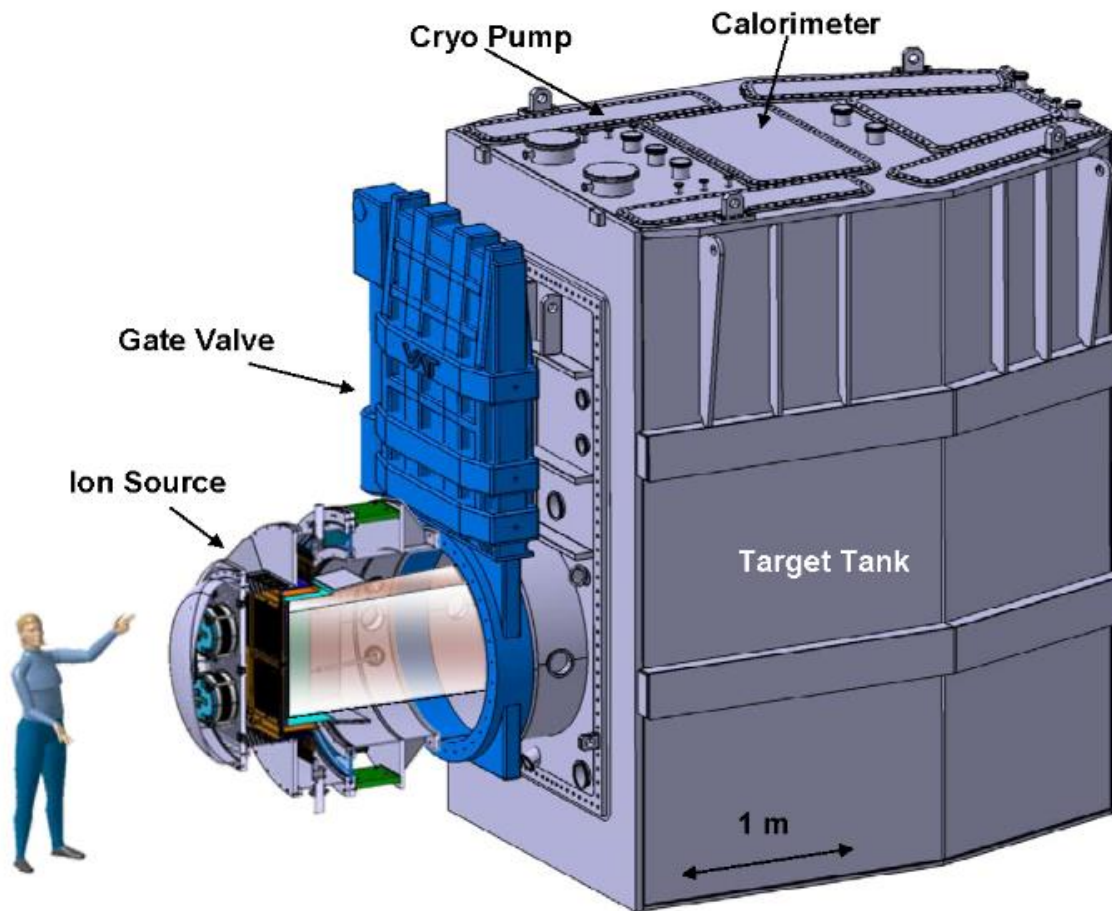


From: www.ipp.mpg.de

Comments

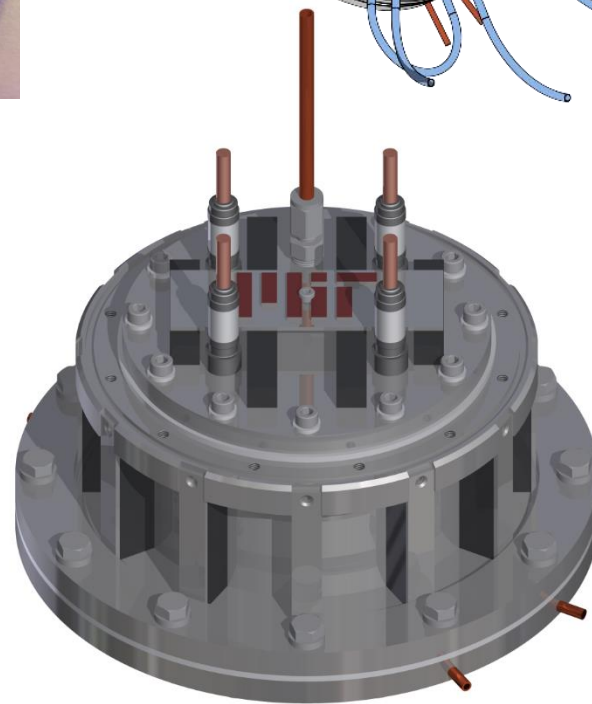
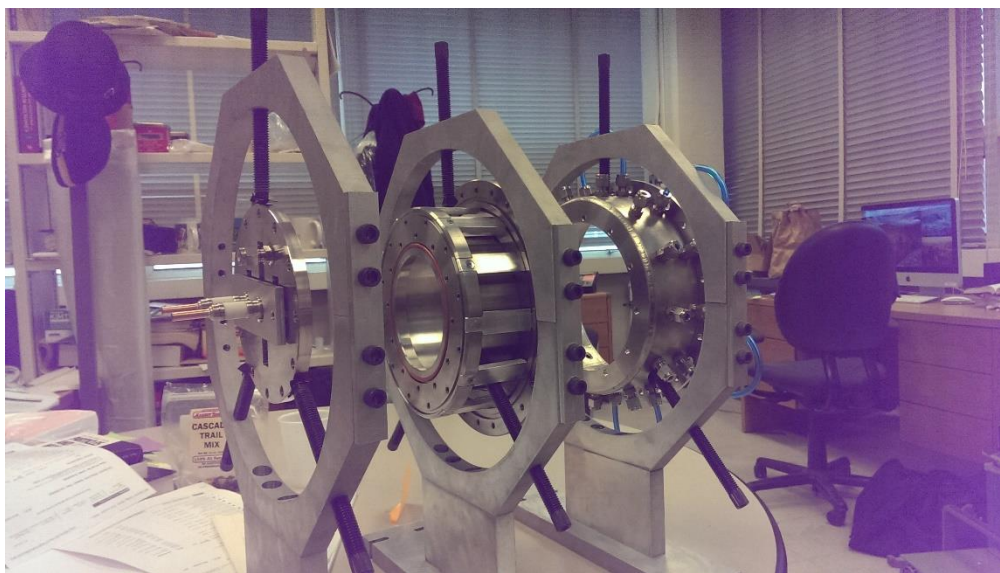
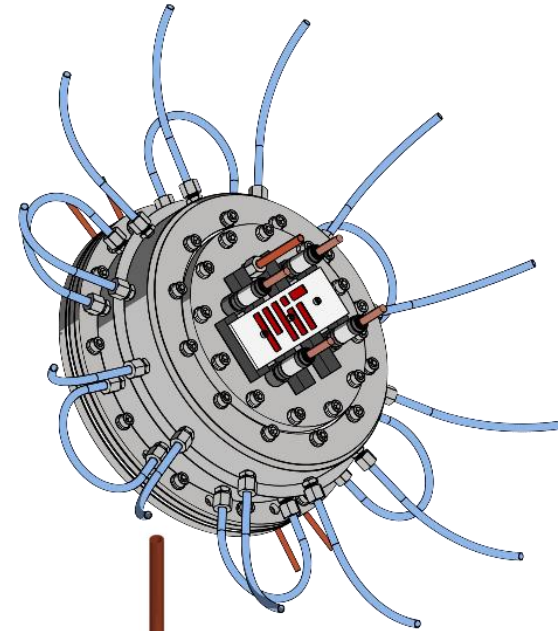
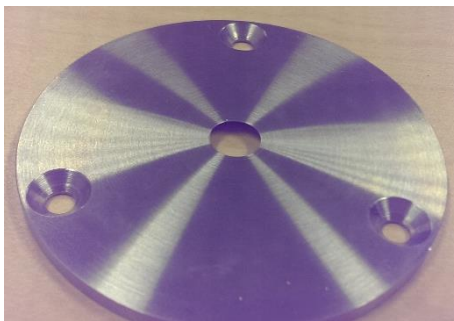
- Large volume plasma production
- Extraction over large area
- Currents in the 10's of A range
- Positive Ion Beams: Sources can produce more current, but neutralization is worse above ~ 100 keV/amu
- Negative Ions: Typical nowadays
- Will talk more about negative ones tomorrow...

ELISE Test Facility for ITER ion source



MIST-1

See other presentation



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<u>RF Heating</u>			
Frequency	5~13 MHz	Somewhat	
Power	few – 100 kW	Yes	
Material/Gas	1e-5 to 1e-3 Torr	Yes	
Extraction	Up to 100 kV	Somewhat	Depends strongly on application!