

# USPAS - Fundamentals of Ion Sources

## 4./5. Beam Quality Parameters I + II

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# Introduction

- A ~1 hour overview on beam quality and how we define various quantities.
- Many topics will be repeated during other parts of the lecture.
- Mostly transversal beam quality.
- Further Reading:
  - Ian Brown – The Physics and Technology of Ion Sources
  - Martin Reiser – Theory and Design of Charged Particle Beams
  - Helmut Wiedemann – Particle Accelerator Physics
  - Many papers... (references will be given on the slides)

# A beam is...

*Ensemble of particles that travel mostly in the same direction (let's use z)*

- Typically:  $v_z \gg v_x, v_y$
- Of course, that's not quite true at the very very beginning, but more later.
- Ion sources:  $k_B T_i \approx \mathcal{O}(\text{eV})$  and extraction voltages  $\mathcal{O}(10\text{kV})$
- Can be comprised of multiple ion species:

$$q_i = Q_i \cdot e, m_i = A \cdot amu \quad (931.5 \text{ MeV}/c^2)$$

# Distributions in 6D Phase Space (+t)

Particle number density:

$$n(x, y, z, p_x, p_y, p_z, t)$$

or

$$n(x, y, z, v_x, v_y, v_z, t)$$

Charge density:  $\rho = q \cdot n$

If the beam is in steady-state (often we extract DC beams from ion sources) one can replace  $t$  with  $z$

Furthermore:  $n(x, x', y, y', z, \Delta p/p)$  (“Trace Space”)

$$x' = \frac{dx}{dz} = \frac{v_x}{v_z}, \quad y' = \frac{dy}{dz} = \frac{v_y}{v_z}$$

# 4D/2D Projections / Slices

If there is no coupling between longitudinal motion and transversal motion the transversal Trace Space density is

$$n(x, x', y, y')$$

Maybe we are even only interested in 2D projections

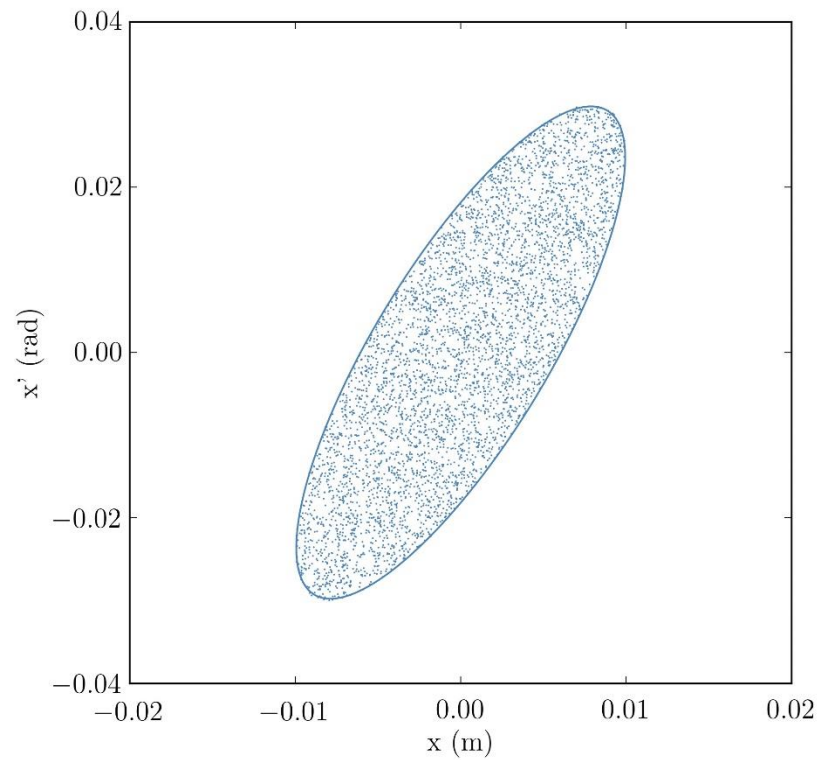
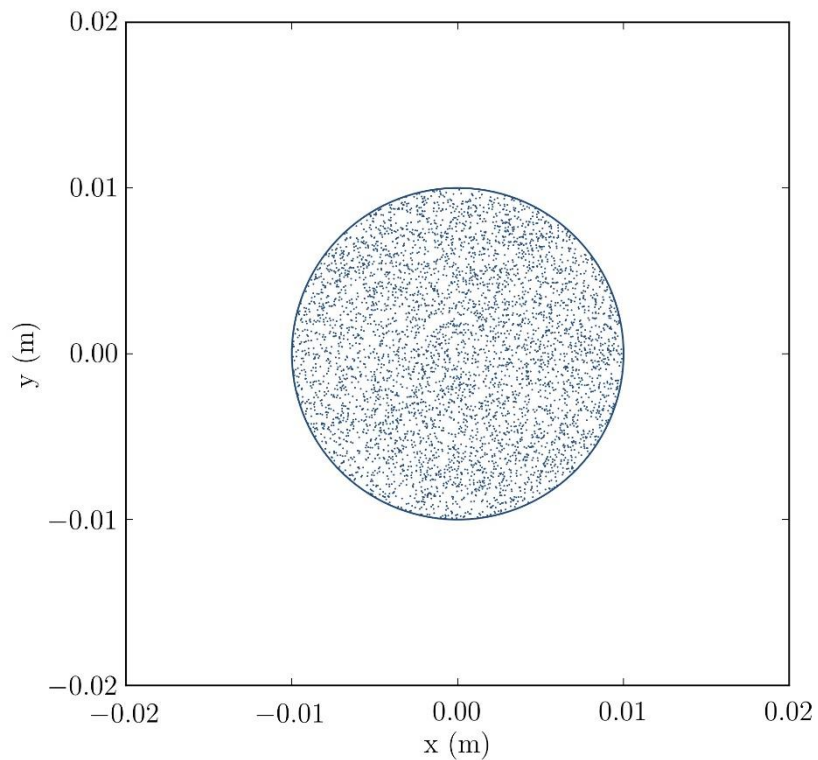
$$n(x, x') = \iint dy dy' n(x, x', y, y')$$

Or slices (interesting in diagnostics and simulations)

$$n(r, r') = n(x, x', y = 0, y' = 0)$$

Because these can tell us something about our beam line transport...

# Trace Space Example



K-V Beam – Projections are uniform ellipses

# Liouville's Theorem

States that for non-interacting particles in a system that can be described by a Hamiltonian, the phase space density is conserved.

$$\frac{dn}{dt} = 0, \text{ or } n = n_0 = \text{const.} \quad \iint d^3q_i d^3P_i = \text{const.}$$

in terms of mechanical momentum

(also true for linear space-charge)

Trace space area:  $A_x = \frac{1}{P} \iint dx dP_x = \frac{1}{\gamma\beta mc} \iint dx dP_x$

# Kapchinsky-Vladimirsky Distribution

The K-V distribution is a uniformly distributed hollow ellipsoid in Trace space:

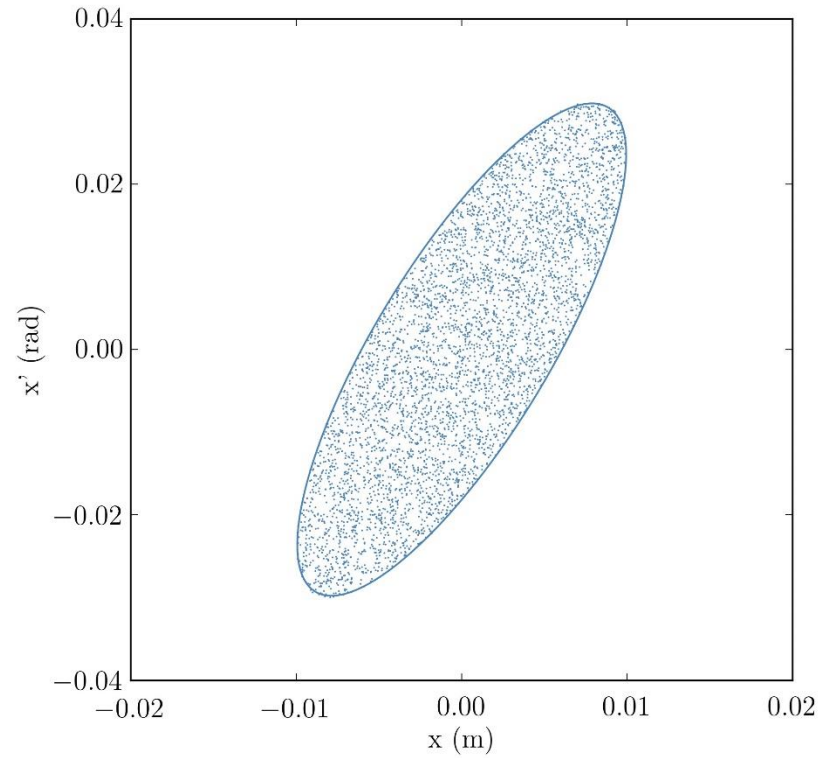
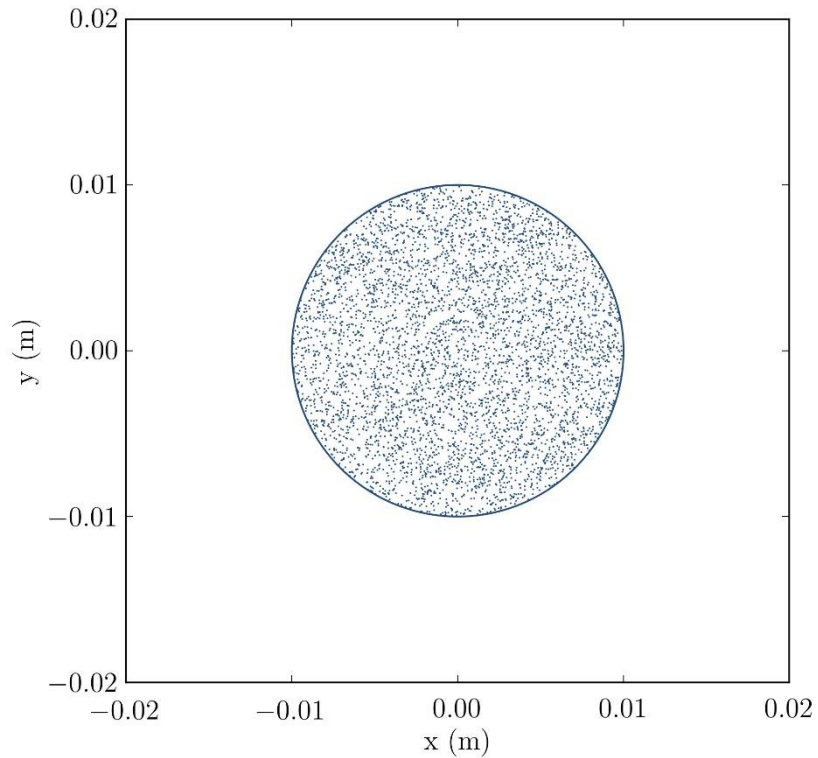
$$f(x, y, x', y') = f_0 \cdot \delta \left( \frac{x_b^2 x'^2 + \sqrt{x_b'^2 x_b^2 - \epsilon_x^2} x x' + x_b'^2 x^2}{\epsilon_x^2} + \frac{y_b^2 y'^2 + \sqrt{y_b'^2 y_b^2 - \epsilon_y^2} y y' + y_b'^2 y^2}{\epsilon_y^2} - 1 \right)$$

with  $x_b, y_b$  the maximum beam extent ( $b$  for 'beam') in  $x$  and  $y$  directions,  $x_b', y_b'$  the maximum angles, and  $\epsilon_x, \epsilon_y$  the (full) beam emittances.

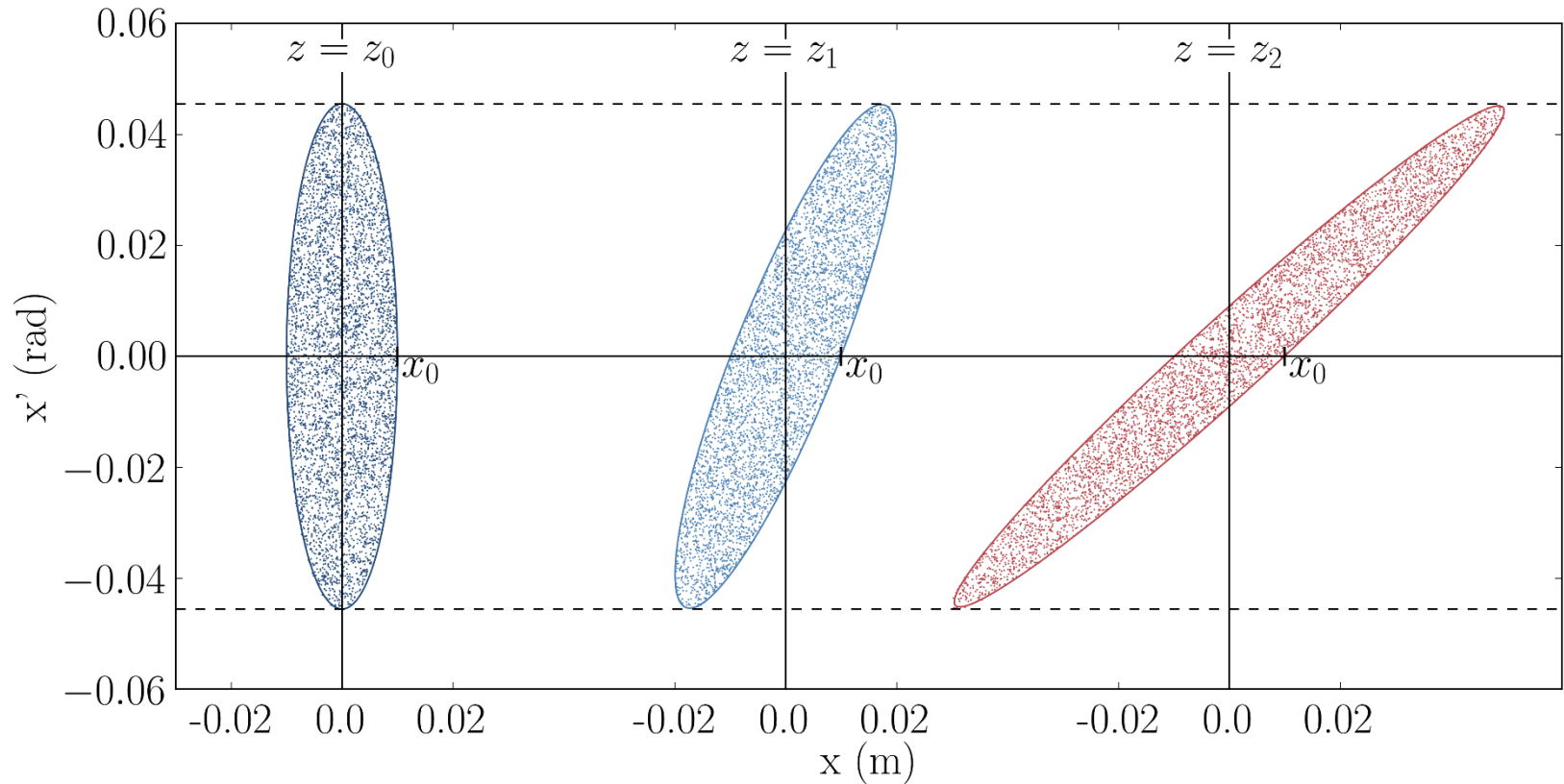
All projections in 2D subspaces are uniformly filled ellipses.



# Trace Space Example



# Phase Space Evolution - Drift



# Geometric Emittance

## Definition from Area

$$\epsilon_x = \frac{A_x}{\pi} \quad [\pi\text{-mm-mrad}]$$

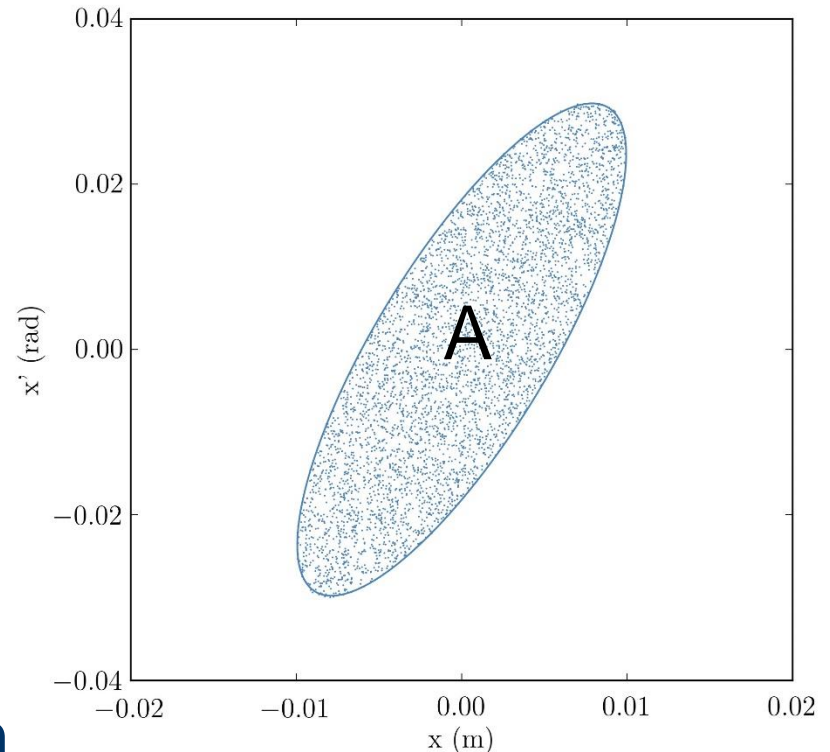
$$A_x = \frac{1}{P} \iint dx dP_x = \frac{1}{\gamma\beta mc} \iint dx dP_x$$

$$A_x = \frac{1}{\gamma\beta} \iint dx dx'$$

## Normalized Emittance:

$$\epsilon_{x,norm.} = \gamma\beta\epsilon_x$$

Const. even under acceleration



# How does Emittance influence Beam Dynamics?

Paraxial equation (similar for y):  $x''(s) + \kappa_x(s)x = 0$

(assuming periodic-focusing and 2 planes of symmetry)

General solution in amplitude-phase notation:

$x(s) = Aw(s)\cos[\Psi(s) + \phi]$  where  $A$  and  $\phi$  are determined by initial conditions

with  $w'' + \kappa w = \frac{1}{w^3}$

Leads to definition of Courant-Snyder Invariant  $\gamma x^2 + 2\alpha x x' + \beta x'^2 = A^2$

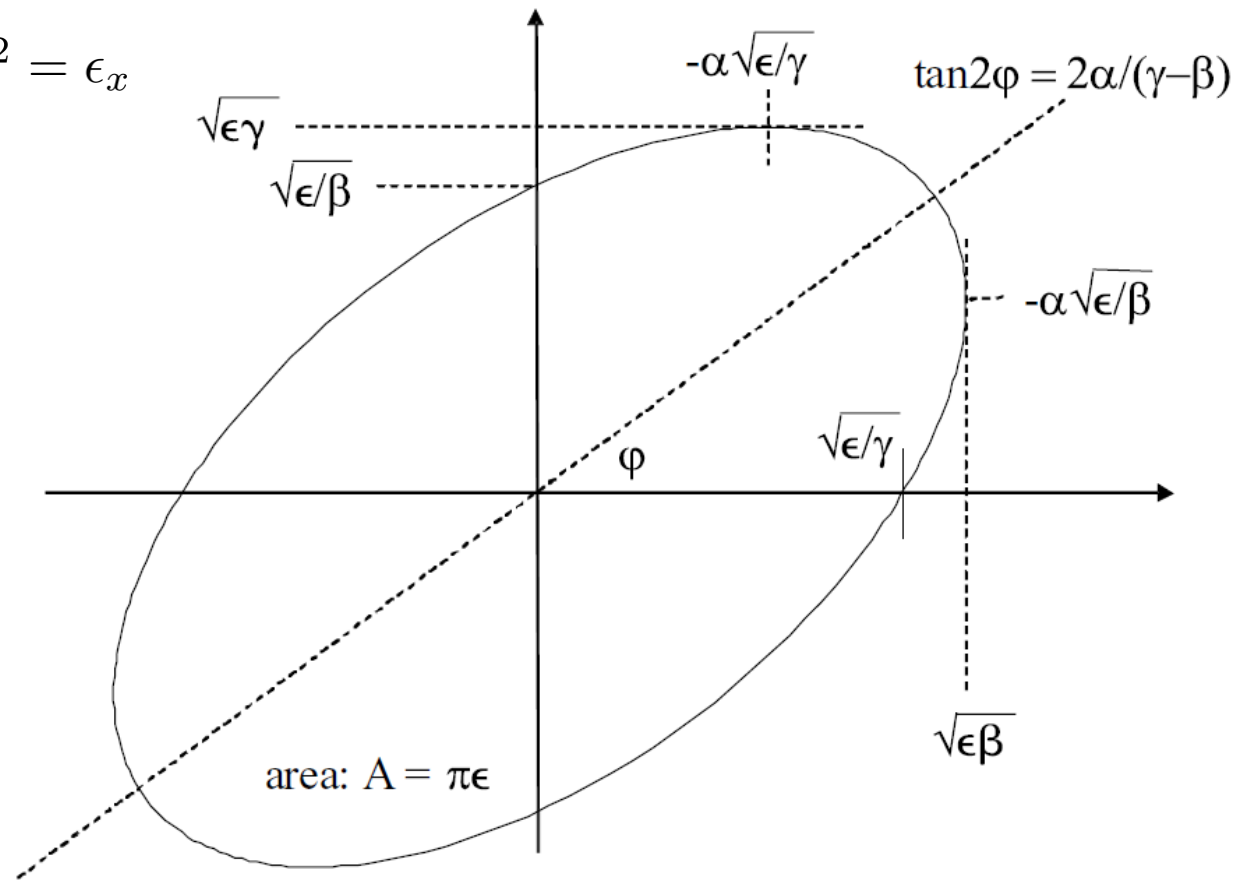
And  $A^2 = \epsilon_x$

(See M. Reiser, *Theory and Design of Charged Particle Beams*)

# Courant-Snyder Invariant and Twiss Parameters

$$\gamma x^2 + 2\alpha x x' + \beta x'^2 = \epsilon_x$$

$$\gamma\beta - \alpha^2 = 1$$

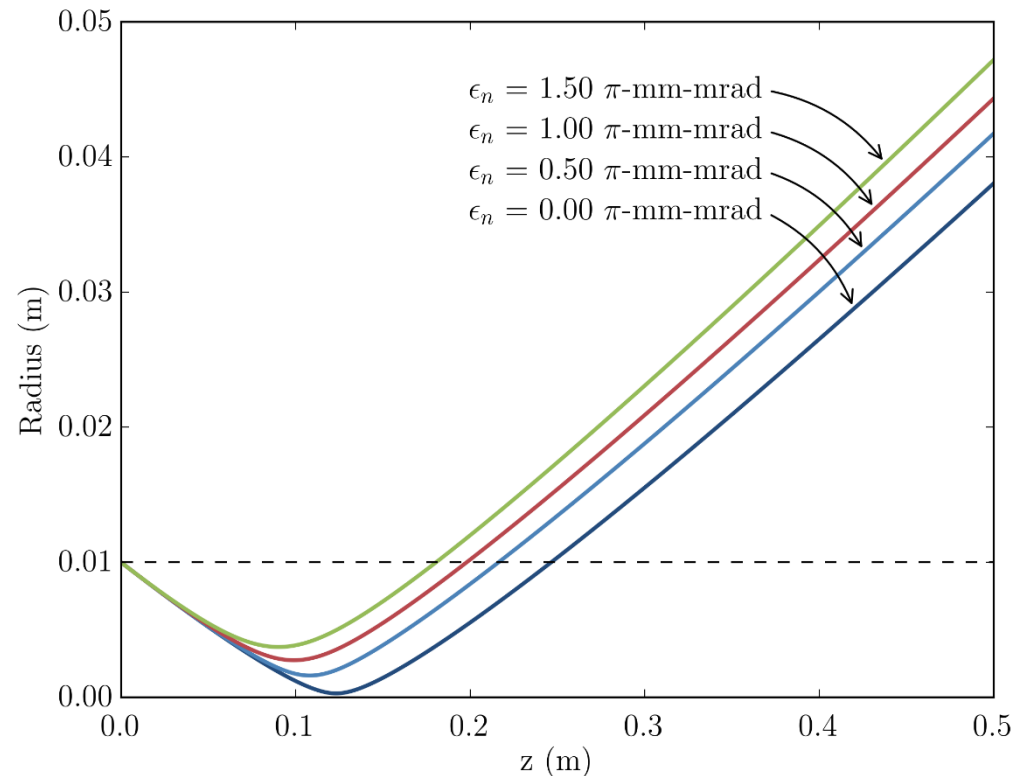


# How does Emittance influence Beam Dynamics?

Courant-Snyder form of envelope equation:

$$x_m'' + \kappa x_m - \frac{\epsilon_x^2}{x_m^3} = 0$$

Emittance works against focusing...

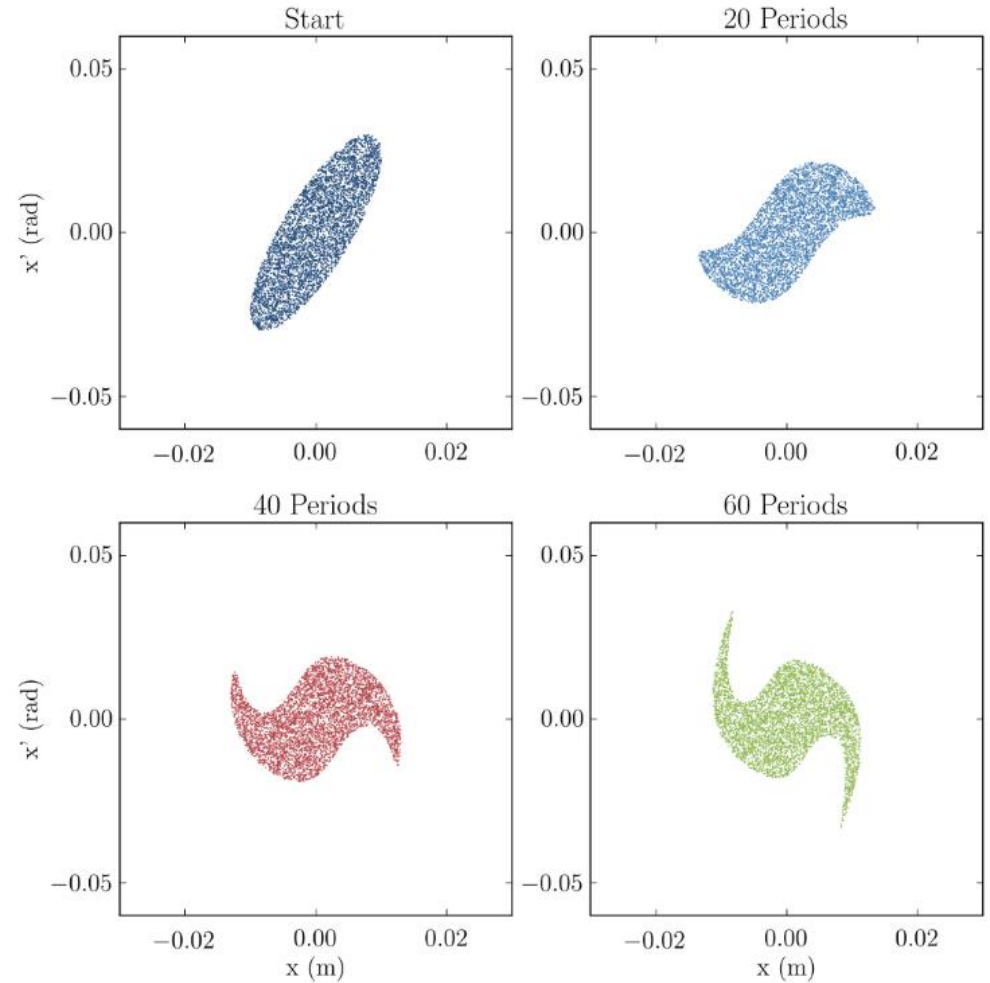


# Why preserve (reduce?) emittance?

- Kind of a no-brainer ;)
- Emittance determines the size of the final focus at a certain focal length from the focusing device.
- Emittance determines the distance beam transport elements have to have.
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- Emittance...the smaller the better...
- And we have a good definition...right?

# Phase Space Evolution – Aberration

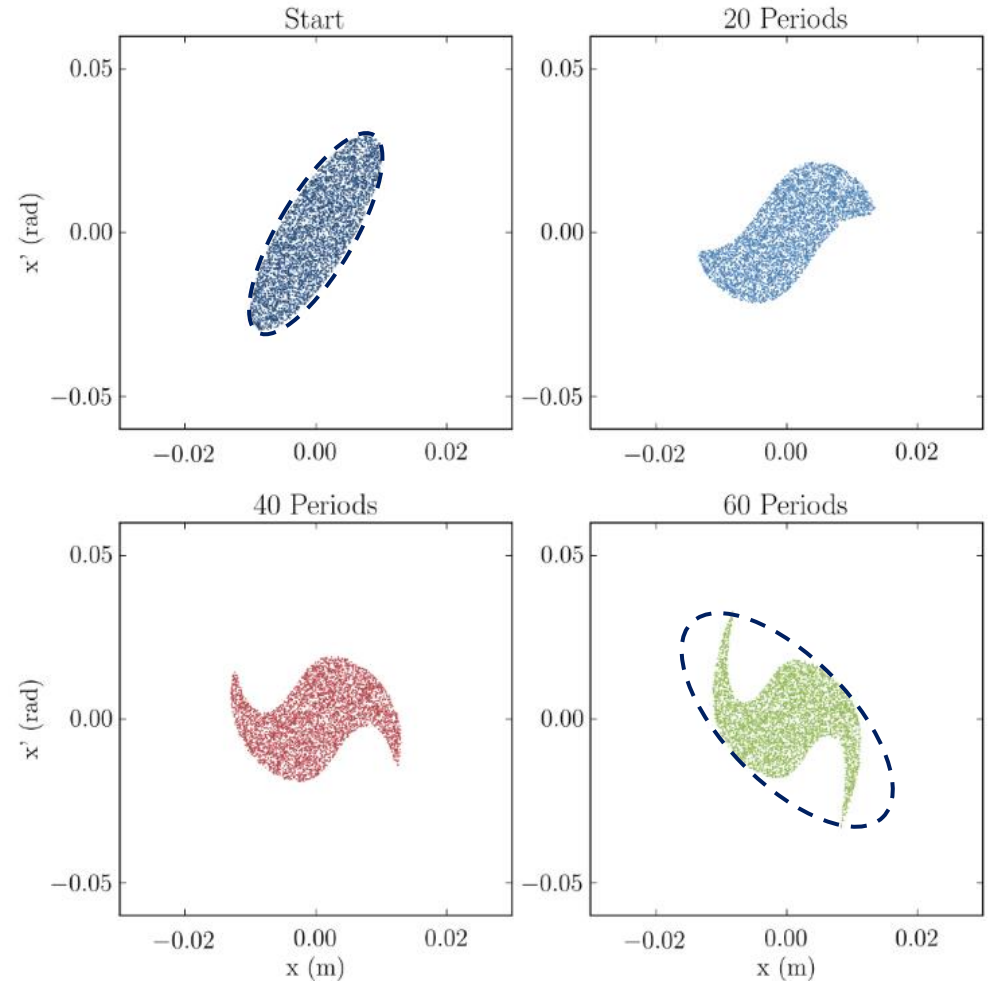
- Simple envelope equation solver with spherical aberration...
- Filamentation of the trace space



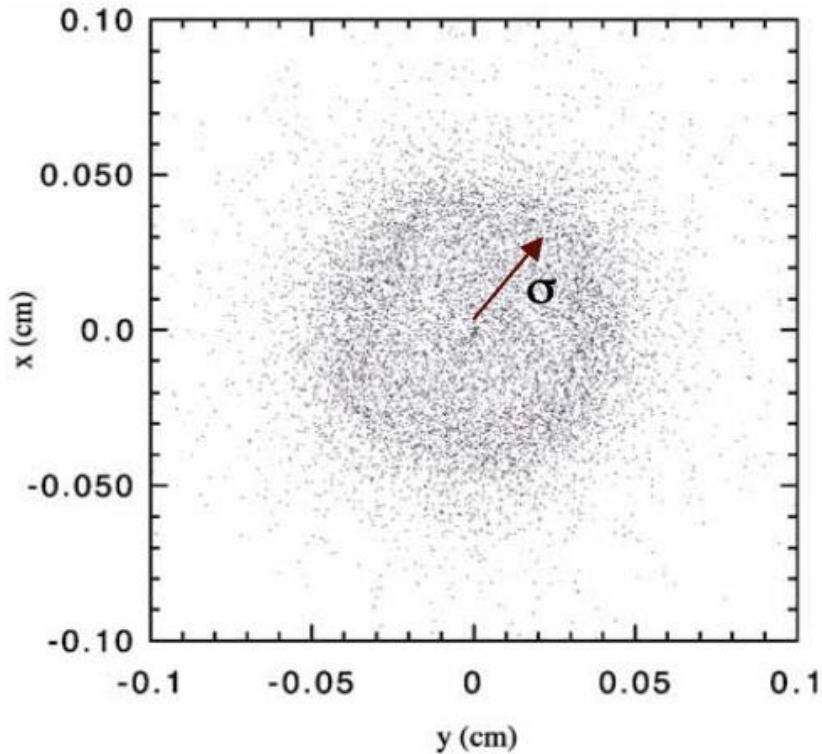


# Phase Space Evolution - Aberration

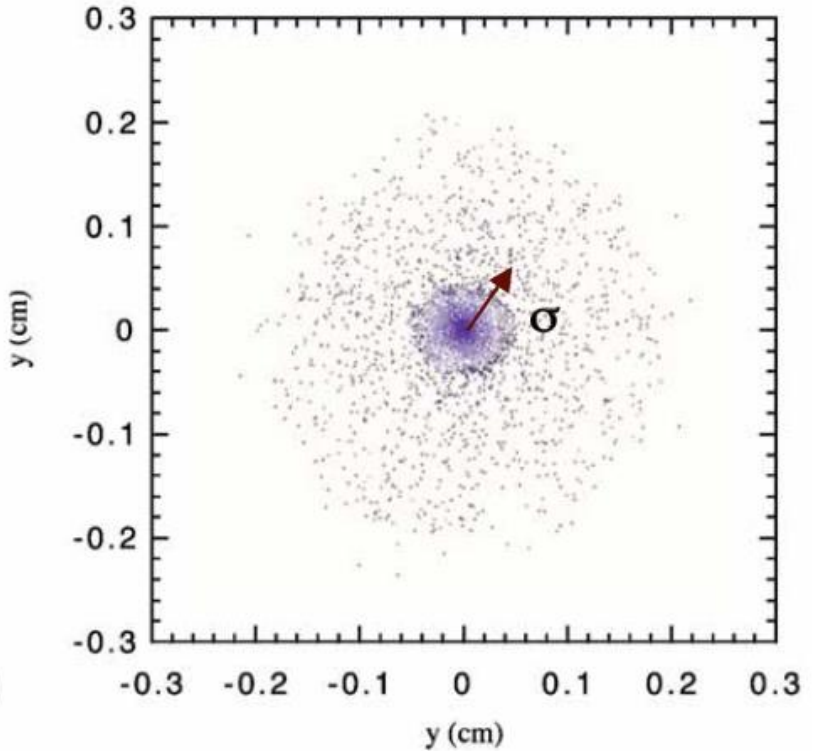
- Simple envelope equation solver with spherical aberration...
- Filamentation of the trace space
- Ellipse surrounding the beam is growing.
- Actual phase space volume is conserved (still Hamiltonian system)



# Other beam Cross-Sections

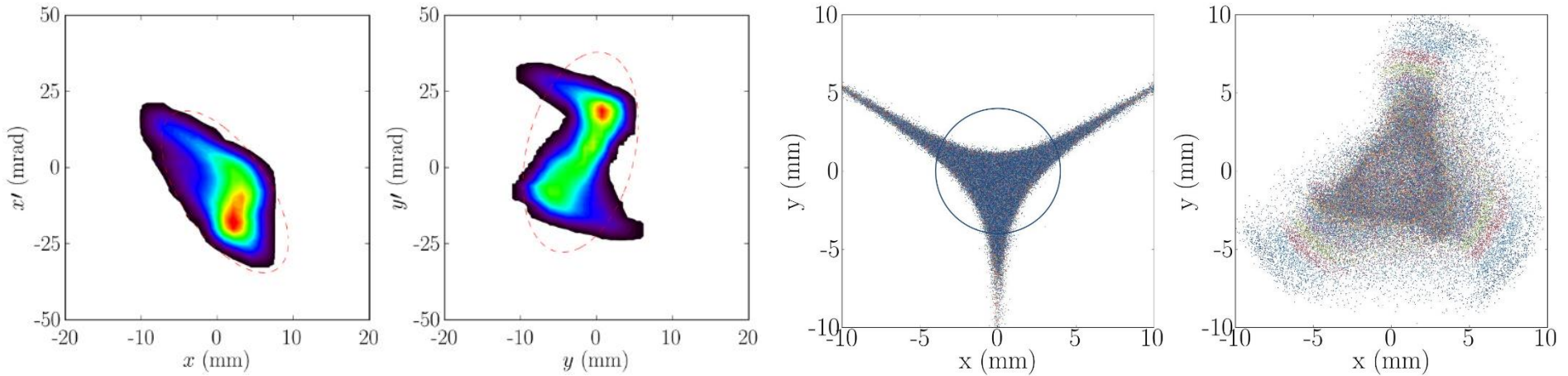


Gaussian beam

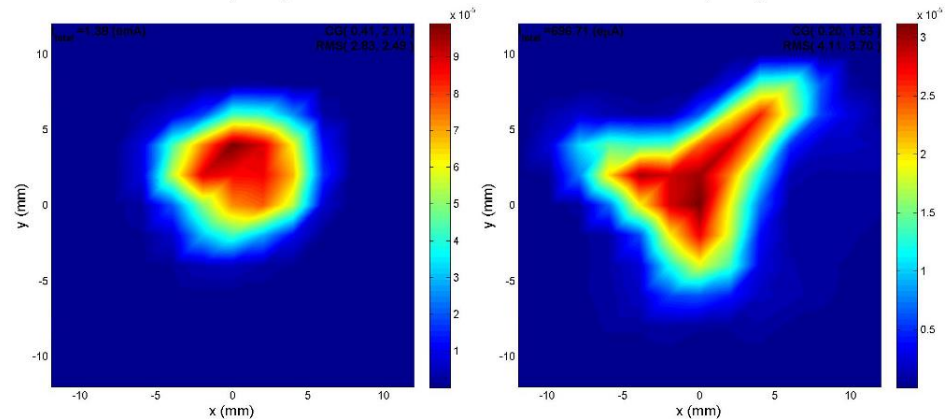


Beam with halo

# Some ECRIS Cross-Sections and Phase Spaces



- Especially in experimental work, we need a more hands-on definition for emittance.



# RMS Emittance

Second moments:

$$\langle x^2 \rangle = \frac{\iiint\!\!\!\int x^2 f(x, y, x', y') dx dy dx' dy'}{\iiint\!\!\!\int f(x, y, x', y') dx dy dx' dy'}$$

$$\langle x'^2 \rangle = \frac{\iiint\!\!\!\int x'^2 f(x, y, x', y') dx dy dx' dy'}{\iiint\!\!\!\int f(x, y, x', y') dx dy dx' dy'}$$

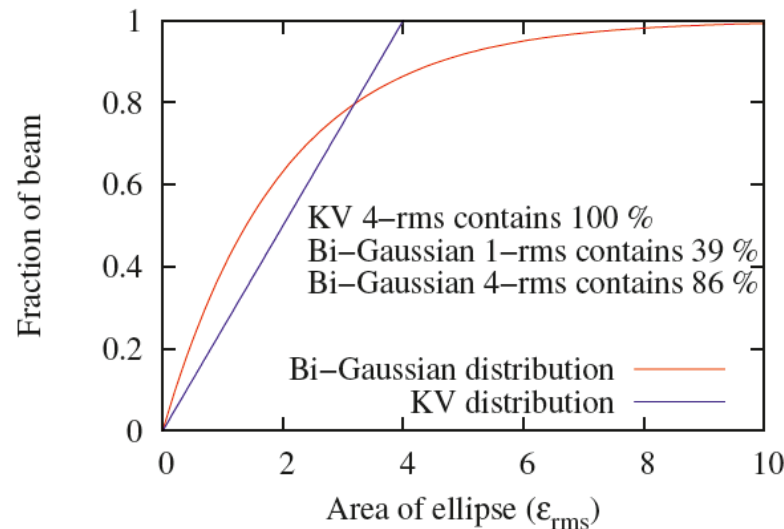
$$\langle xx' \rangle = \frac{\iiint\!\!\!\int xx' f(x, y, x', y') dx dy dx' dy'}{\iiint\!\!\!\int f(x, y, x', y') dx dy dx' dy'}$$

$$\epsilon_{rms} = \sqrt{\langle x^2 \rangle \langle x'^2 \rangle - \langle xx' \rangle^2} \quad [mm\text{-}mrad]$$

# How does this compare to full emittance?

Well, that depends...on the actual distribution.

- K-V Distribution:  $\epsilon_x = 4\epsilon_{x,rms}$
- Waterbag Distribution:  $\epsilon_x = 6\epsilon_{x,rms}$
- (Bi-)Gaussian Distribution:  $\epsilon_x = n^2\epsilon_{x,rms}$  if truncated at  $n \cdot \sigma$



# Effective Emittance/Beam Size

- For analytical calculations (e.g. using the K-V distribution) It can be useful to define  $\epsilon_x = 4\epsilon_{x,rms}$  as “effective emittance” as it corresponds to the K-V distribution which is often used for first order analytical investigations.
- Many authors have adopted this method.
- Similar:  $X = 2\tilde{x} = s (\bar{x^2})^{1/2}$  the “effective beam radius”

# Equivalent Emittance

Often, what we are interested in is really how much beam can we transport along a beamline, through a series of apertures, etc. so it makes sense to define a 90% emittance. (or look at 90% versus maximum to determine halo)

→ The emittance contour within which 90% of the beam sits. (Often this is close to 4-rms).

$$\epsilon_{x-90\%}$$

# Brightness

- The brightness is commonly defined as current density per unit solid angle.

$$B = \frac{J}{d\Omega} = \frac{dI}{dSd\Omega}$$

- Or in terms of the transversal projections

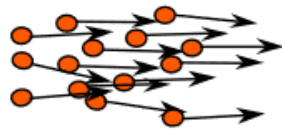
$$\bar{B} = \frac{2I}{\pi^2 \epsilon_x \epsilon_y} \left[ \frac{A}{m^2 \cdot rad^2} \right]$$

$$B_n = \frac{B}{\beta^2 \gamma^2} \quad B_{90\%} = \frac{2 \cdot 0.9 \cdot I}{\pi^2 \epsilon_{x-90\%} \epsilon_{y-90\%}}$$



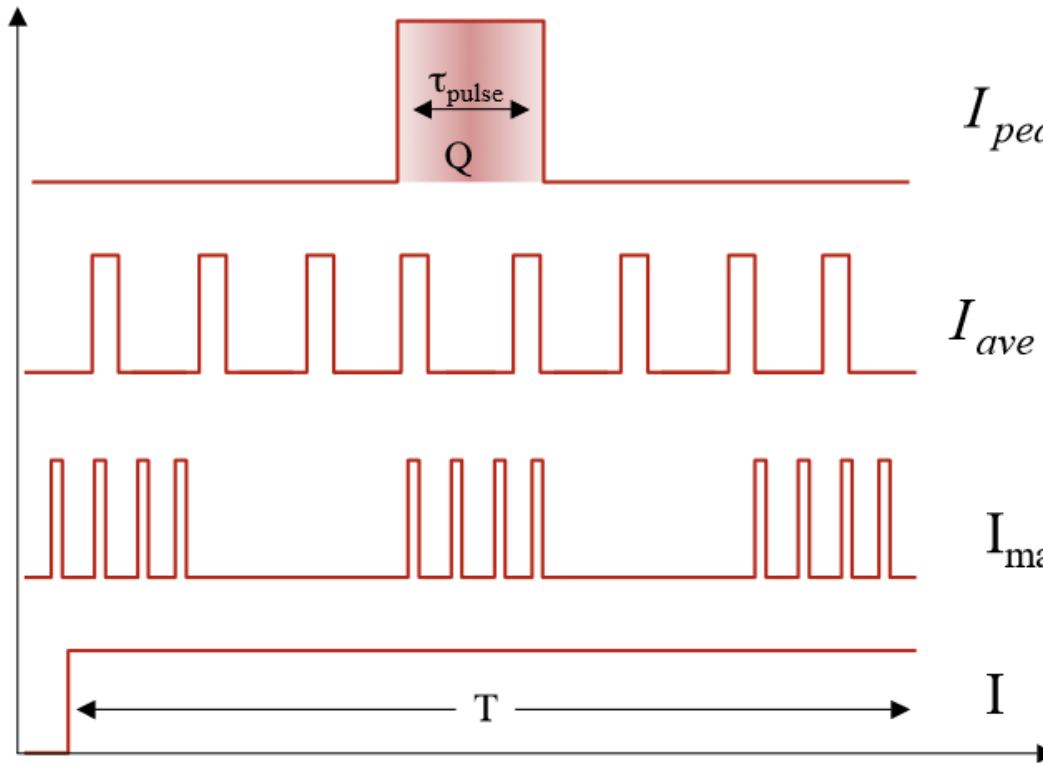
# Longitudinal Beam Properties

## Bunch Currents



$$I \sim ne\langle v_z \rangle$$

$$\text{Duty factor} = \frac{\sum \tau_{pulse}}{T}$$



$$I_{peak} = \frac{Q}{\tau_{pulse}}$$

$$I_{ave} = \frac{Q_{tot}}{T}$$

$I_{macro}$

$I$

# Longitudinal Beam Properties

## Phase Space Variables

Relative energy variation:  $\delta = \frac{\Delta E}{E_0} = \frac{E - E_0}{E_0}$

Relative time:  $\tau = t - t_0$

Momentum variation:  $\frac{\Delta \gamma}{\gamma}, \frac{\Delta p}{p}$

Path along beam line:  $z$  or  $s$

# Longitudinal Beam Properties

## Comments

Energy spread in beam comes from:

- Source HV potential stability
- Ions are created at different potentials inside the sheath
- Plasma instabilities
- Temperature  $kT$

Measure energy spread of low energy beams? E.g. Retarding Field Analyzer.

# What reduces beam quality in LEBT?

- Source performance and plasma parameters.
- Beam mismatch: Large beam going through solenoid/ Einzel lens/Dipole, will have aberrations.
- Coupling of longitudinal and transversal motion. (e.g. Energy spread going through a dipole magnet)
- Space Charge: Repulsive force of beam ions toward each other.
- Over-focusing will increase space charge density locally.
- Transporting unwanted ion species/electrons.

# How can we preserve beam quality?

- Well designed beam lines.
- Stable power supplies.
- Separation of ions early on.
- Correct beam line pressure for highly charged ions, fragile ions and protons. (Space Charge Compensation)

# Summary of Beam Quality Parameters

Current

Divergence

Emittance

Brightness

Energy Spread

Purity

Some have many different definitions and uses...